

# Elliptic flow of D mesons in Pb-Pb collision for ALICE

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# Introduction

We want to verify the possibility at ALICE to discriminate between different elliptic flow models.

So we want to develop a method to introduce an elliptic flow in a (isotropic) Montecarlo Generator. Characteristics:

- has to be very fast
- has to use different values for  $v_2$
- manages D mesons differently from other particles

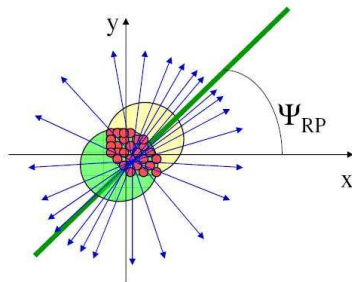
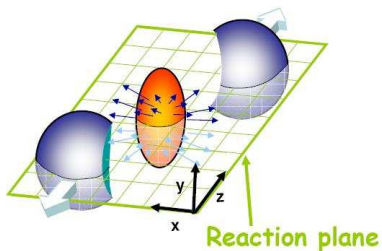
# Introduction (2)

2 ways to generate elliptic flow:

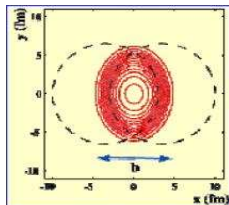
- 1 Build up a full hydrodynamical simulation to generate elliptic flow.
  - More precise, possible predictions
  - takes a lot of time and work
- 2 Generates expected particles from standard generators (PYTHIA+HIJING). Then change particles momenta distribution to introduce elliptic flow.

# The Elliptic Flow (1)

When 2 nucleus collides with impact parameter  $b \neq 0$  at high energy, interaction region is not isotropic. Thus, if the system thermalize, there are different pressures on x and y axis that force an anisotropy on momenta distribution.



## The Elliptic Flow (2)



If we expand with Fourier the azimuthal particles distribution we get

$$\frac{d^3N}{p_t dp_t dy d\phi} = \frac{1}{2\pi} \frac{d^2N}{p_t dp_t dy} (1 + 2v_1 \cos(\phi - \Psi) + 2v_2 \cos[2(\phi - \Psi)] + \dots)$$

- 0<sup>th</sup> : Radial flow
- 1<sup>st</sup> : Direct flow
- 2<sup>nd</sup> : Elliptic flow

$$v_n = \frac{\int d\phi \cos(n\phi) \frac{dN}{d\phi}}{\int d\phi \frac{dN}{d\phi}} = \langle \cos(n\phi) \rangle$$

## The Elliptic Flow (3)

At midrapidity, for geometrical reasons, all the terms  $v_n$  with odd  $n$  are very low.

So if it's true that  $v_2 \gg v_4 \gg v_6$  we get

$$\frac{2\pi}{N_0} \frac{dN}{d\phi} = 1 + 2v_2 \cos[2(\phi - \Psi)]$$

where  $v_2$  is called *elliptic flow* and is possible to calculate it as mean value of  $\cos[2(\phi - \Psi)]$

# Input of $v_2$ (1)

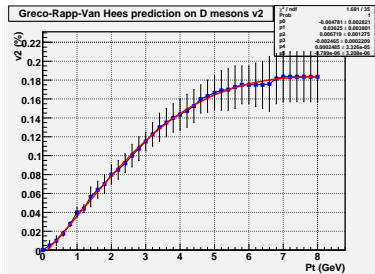
A lot of different models of  $v_2$  are going to be tested at LHC. It has to be very easy to add any new trend of  $v_2$  in the tool.

I'm working on:

- Langevin hydrodynamics. (V. Greco, H. Van Hees, R. Rapp)
- Ideal hydrodynamic models. (J. Y. Ollitrault)
- Covariant transport model. (D. Molnar)
- CGC + hydro + cascade models. (T. Hirano, M. Nardi)
- agnostic extrapolation of  $v_2$  from RHIC points.
- a model based on angular momentum conservation.  
(Becattini)

# Models of $v_2$ (1) - Langevin

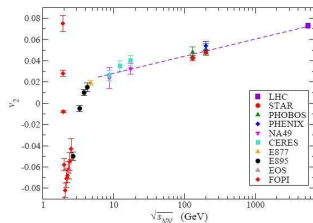
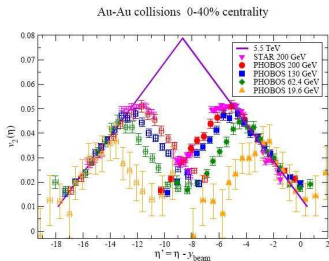
The default value of  $v_2$  used for D mesons is taken from the Langevin model with a coalescence-fragmentation hadronization.





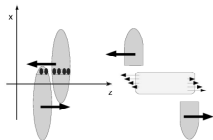
## Modelli di $v_2$ (2) - Extrapolations from RHIC

There are some universal trends (scaling dof  $N_{ch}/N_{ch}^{pp}$  with  $\langle N_{part} \rangle$ , scaling of  $\langle N_{ch} \rangle / \langle N_{part} \rangle$  with  $\ln^2 \sqrt{s}$ ) that allow us to build expected pseudorapidity distribution. We can use it to extend  $v_2(\eta)$  from RHIC to LHC.



This way pions are expected to have an elliptic flow between 7% and 8%.

## Models of $v_2$ (3) - Angular momentum



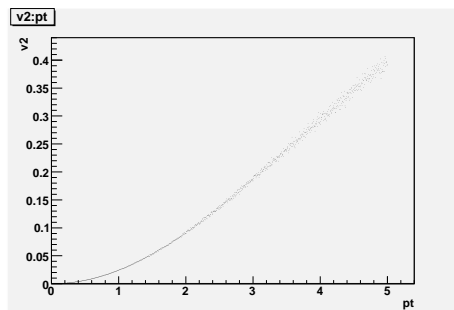
Bjorken hypothesis implies  $\frac{\partial v_z}{\partial x_i} \approx 0$  but these hypotheses doesn't allow angular momentum global conservation. If we don't make it we find out that

$$\epsilon_0 \gamma_0^3 \frac{\partial u_i}{\partial t} \Big|_{t=0} = -\frac{1}{4} \frac{\partial \epsilon \gamma^2}{\partial x_i} \Big|_{t=0} + \frac{1}{4} 2 \epsilon_0 \gamma_0^4 v_{z0} \frac{\partial v_{z0}}{\partial x_i} \Big|_{t=0}$$

- First term: spatial anisotropy (standard)
- Second term: new term due to angular momentum conservation

## Models of $v_2$ (4) - Angular momentum

It's not trivial getting from the equation before to an analytical expression for  $v_2$ . If we use a spinning rigid sphere (not realistic) we have this result

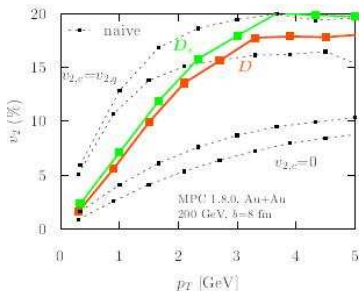


## Models of $v_2$ (5) - Covariant transport

The system is made of 3 massless quarks ( $u, d, s$ ) and 1 massive ( $c$ ). It evolves through 2 body into 2 body collisions (both elastic and inelastic).

Hadronization is made through coalescence.

Not all the processes are taken into account  $\Rightarrow$  Overestimates x-sections.



# Test values

in order to be able to test the program, we're at the moment using the following settings for  $v_2$ :

- a constant  $v_2$  of 0.07 for kaons and pions (expected midrapidity value).
- a constant  $v_2$  of 0.05 for protons.
- a  $p_T$  dependent value of  $v_2$  for D mesons based on a coalescence-fragmentation scheme (Van Hees - Greco - Rapp).

# The tool

The tool works in sequential steps:

- 1 Generation of the particles via cocktail (HIJING + PYTHIA) of generators.
  - Heavy quarks switched off in HIJING, they are generated by PYTHIA. The number of  $c\bar{c}$  couples reproduces pQCD next to leading order calculations for Pb-Pb (MNR).
  - The generation is isotropic.
- 2 Reading of the Kinematics: are selected pions, kaons, protons and D mesons.
  - We throw away particles coming from D decays.
  - Kinematic cuts (pseudorapidity, decay region...)
- 3 Introduction of the elliptic flow  $\Leftarrow$
- 4 Decay of the D mesons  $\Leftarrow$

# Introduction of the elliptic flow (1)

The standard method (used also in the afterburner class of Aliroot) to introduce elliptic flow in a set of isotropically generated particles is (Poskanzer - Voloshin)

arXiv:nucl-ex/9805001v2 24 Jun 1998  
Methods for analyzing anisotropic flow in relativistic nuclear collisions  
A. M. Poskanzer<sup>1)</sup> and S. A. Voloshin<sup>2)</sup>

## V. SIMPLE WAY TO INTRODUCE FLOW IN A MONTE-CARLO EVENT GENERATOR

Sometimes in order to investigate different detector effects or the reliability of the method, one needs to introduce anisotropic flow into a Monte-Carlo event generator. It can be done by changing the azimuthal angle of each particle (and consequently changing the density in the azimuthal angle space)

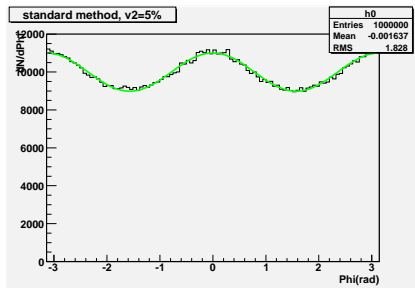
$$\phi \rightarrow \phi' = \phi + \Delta\phi, \quad (34)$$

where

$$\Delta\phi = \sum_n \frac{-2}{n} v_n \sin(n(\phi - \psi_0)), \quad (35)$$

## Introduction of the elliptic flow (2)

for  $v_2$  higher than 5% our distribution is not following the expected  $\frac{dN}{d\phi'} = 1 + 2v_2 \cos(2\phi')$  (Fourier expansion).

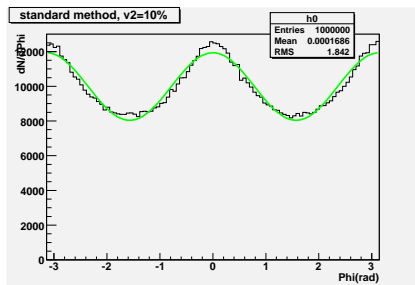


green: expected trend and higher harmonics; black: simulation result.



## Introduction of the elliptic flow (2)

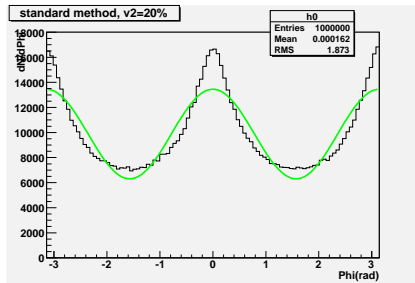
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green: expected trend and higher harmonics; black: simulation result.

## Introduction of the elliptic flow (3)

We start from isotropic distributions  $\Rightarrow \frac{dN}{d\phi} = K$  and we want to get a distribution like  $\frac{dN}{d\phi'} = 1 + 2v_2 \cos(2\phi')$

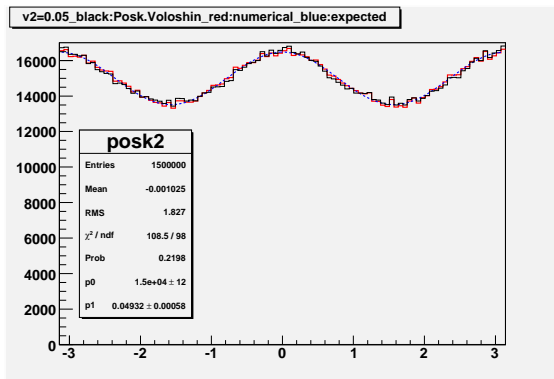
$$K = \frac{dN}{d\phi} = \frac{dN}{d\phi'} \frac{d\phi'}{d\phi} = [1 + 2v_2 \cos(2\phi')] \frac{d\phi'}{d\phi}$$

so, very easily

$$\phi' = \phi - v_2 \sin(2\phi')$$

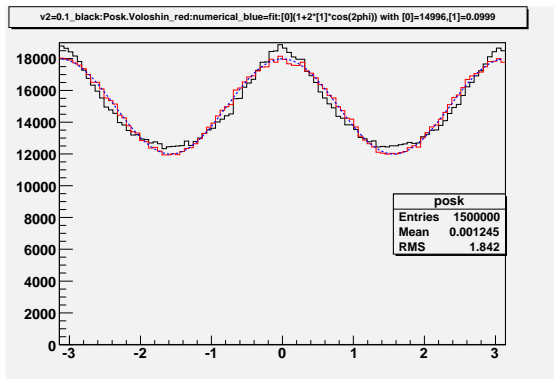
which is the same of the Poskanzer-Voloshin equation  $\phi' = \phi - v_2 \sin(2\phi)$  only in the approximation of low  $v_2$  ( $\phi' \sim \phi$ ). we can solve the correct equation above using Newton roots method, and we get...

# Introduction of the elliptic flow (4)



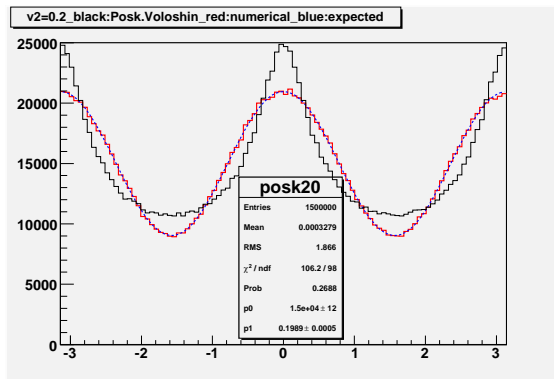
Blue: expected trend; Red: Simulation result; Black: Poskanzer-Voloshin simulation result.

# Introduction of the elliptic flow (4)



Blue: expected trend; Red: Simulation result; Black:  
Poskanzer-Voloshin simulation result.

# Introduction of the elliptic flow (4)

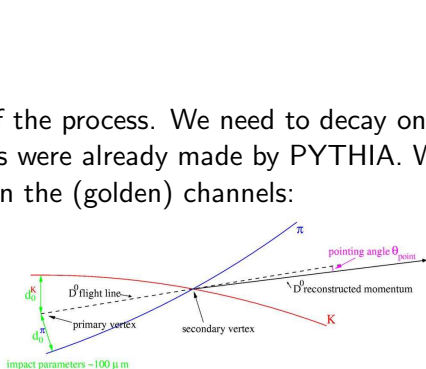


Blue: expected trend; Red: Simulation result; Black: Poskanzer-Voloshin simulation result.

# Decayer (1)

The decayer is the last step of the process. We need to decay only D mesons, all the other decays were already made by PYTHIA. We force the D mesons to decay in the (golden) channels:

- $D^0 \rightarrow K^+ \pi^-$
- $D^+ \rightarrow K^- \pi^+ \pi^+$
- $D_s^+ \rightarrow K^- K^+ \pi^+$



## Decayer (2)

The 2-body decay is made generating a isotropic decay in the CM rest frame and then boosting in the lab frame.

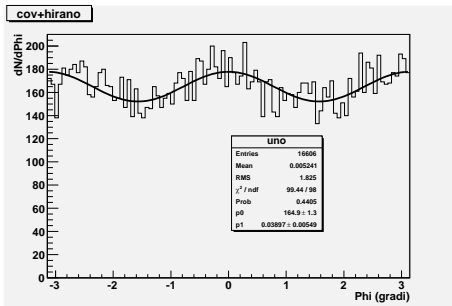
The 3-body decay is made using 2 subsequential 2-body decay ( $D \rightarrow 1 + 23 \rightarrow 2 + 3$ ).

Matrix element corrections and phase space weights are calculated as in PYTHIA decayer.

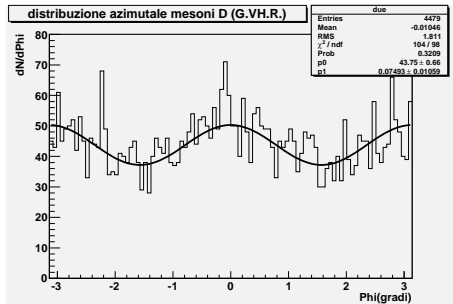
Resonance decays are taken in account with their relative probability.



# Results (1)

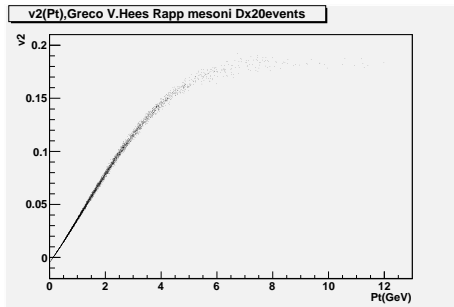


Pions azimuthal distribution in 1 mid-central Pb-Pb event (CGC+hydro for  $p_t < 3$  GeV, covariant transport elsewhere)

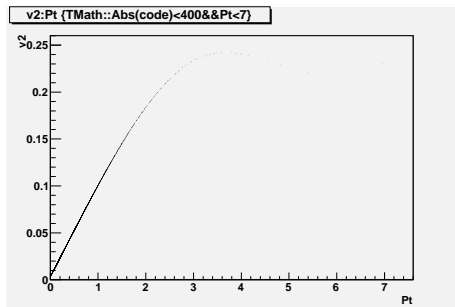


D mesons azimuthal distribution in 20 mid-central Pb-Pb events. (Langevin hydro)

## Results (2)

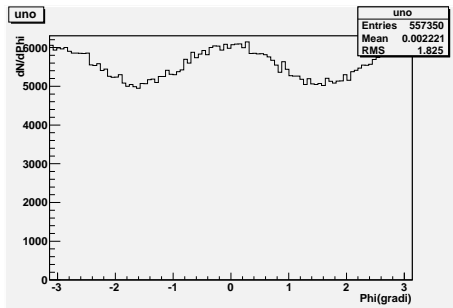


$v_2(p_t)$  for D mesons in 20  
smidcentral Pb-Pb events (Langevin  
hydro)

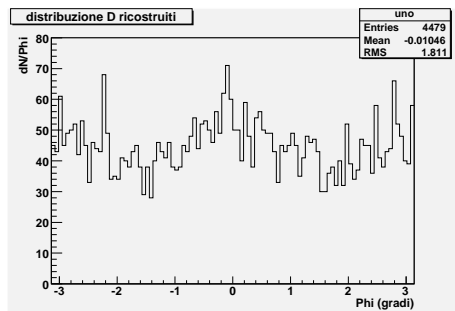


$v_2(p_t)$  for pions in 1 midcentral  
Pb-Pb event (covariant transport)

# Results (3)



azimuthal distribution of particles  
after decayer



azimuthal distribution of  
reconstructed D mesons

# Conclusions

Is this method fast? Does it use too much memory?

- Time to run a 1 event Pb-Pb generation (PYTHIA + HIJING, ITS switched off) on standard machine:  $\approx$  5 minutes
- Time to run afterburner and decayer on standard machine:  $\approx$  2 minutes and a half.
- Space on disk used by all the files created by the generator: 6,3 Mb
- Space on disk used by the files created by this tool: 8 Mb