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Stefano Gariazzo

IFIC, Valencia (ES)
CSIC – Universitat de Valencia

gariazzo@ific.uv.es

<http://ificio.uv.es/~gariazzo/>

Dark radiation candidates: light sterile neutrinos and thermal axions

COSMOS meeting, Ferrara (IT), 26–27/06/2018

1 *Light sterile neutrinos*

- Why a sterile neutrino
- Cosmological constraints

2 *Thermal axions*

- Motivation and theory
- Cosmological constraints

3 *Conclusions*

Neutrino Oscillations

Analogous to CKM mixing for quarks:

[Pontecorvo, 1958]

[Maki, Nakagawa, Sakata, 1962]

$$\nu_\alpha = \sum_{k=1}^3 U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau)$$

ν_α flavour eigenstates, $U_{\alpha k}$ PMNS mixing matrix, ν_k mass eigenstates.

Current knowledge of the 3 active ν mixing: [de Salas et al. (2018)]

$$\Delta m_{ji}^2 = m_j^2 - m_i^2, \theta_{ij} \text{ mixing angles}$$

NO: Normal Ordering, $m_1 < m_2 < m_3$

IO: Inverted Ordering, $m_3 < m_1 < m_2$

$$\Delta m_{21}^2 = (7.55^{+0.20}_{-0.16}) \cdot 10^{-5} \text{ eV}^2$$

$$|\Delta m_{31}^2| = (2.50 \pm 0.03) \cdot 10^{-3} \text{ eV}^2 \text{ (NO)}$$
$$= (2.42^{+0.03}_{-0.04}) \cdot 10^{-3} \text{ eV}^2 \text{ (IO)}$$

$$\sin^2(\theta_{12}) = 0.320^{+0.020}_{-0.016}$$

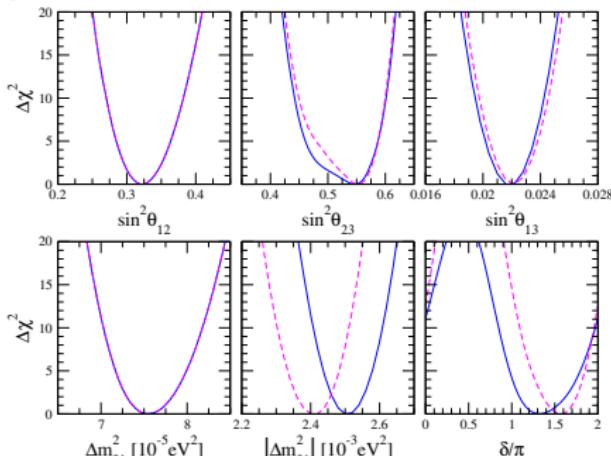
$$\sin^2(\theta_{13}) = 0.0216^{+0.008}_{-0.007} \text{ (NO)}$$

$$= 0.0222^{+0.007}_{-0.008} \text{ (IO)}$$

$$\sin^2(\theta_{23}) = 0.547^{+0.020}_{-0.030} \text{ (NO)}$$

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First hints for $\delta_{CP} \simeq 3/2\pi$



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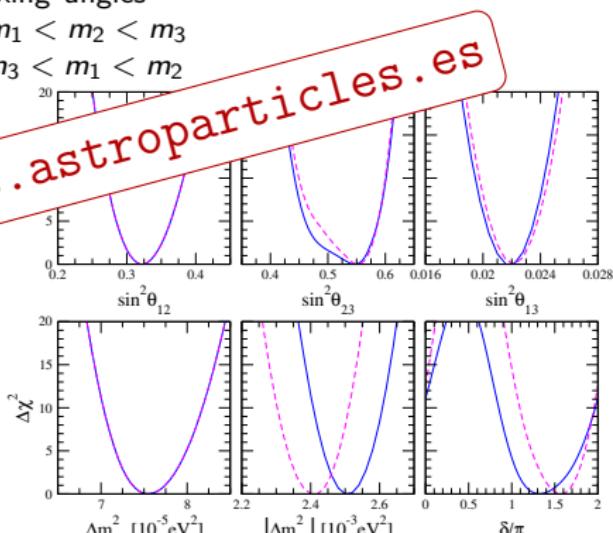
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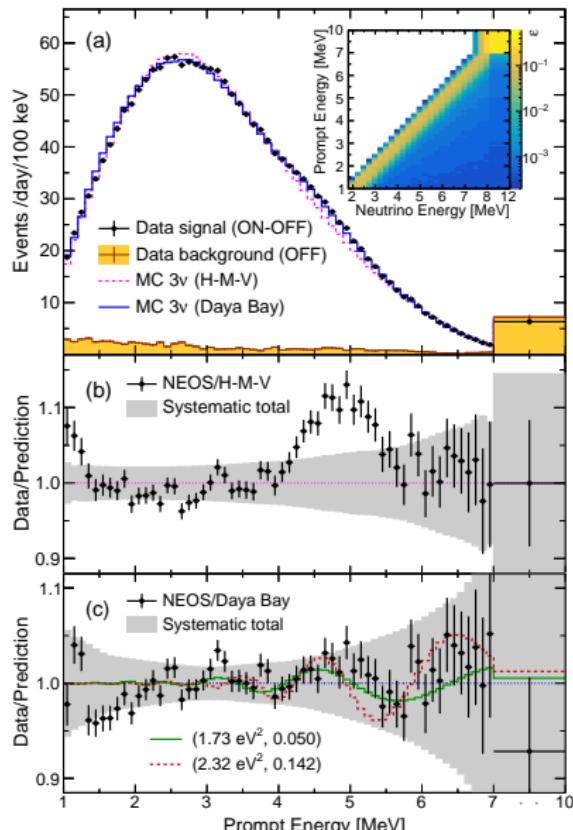
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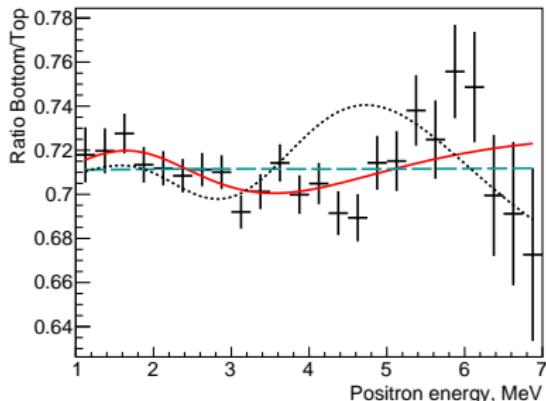


Recent results...

[NEOS, PRL 118 (2017) 121802]



[DANSS, arxiv:1804.04046]

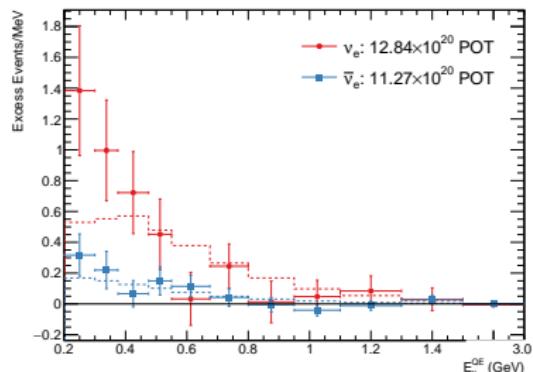
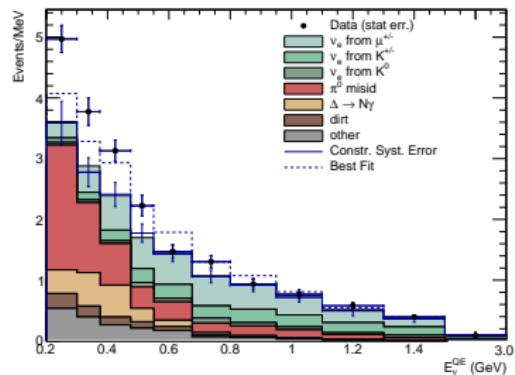


first *model independent*
indications in favor
of SBL oscillations

DANSS alone gives a
 $\Delta\chi^2 \simeq 13$ in favor of
a light sterile neutrino!

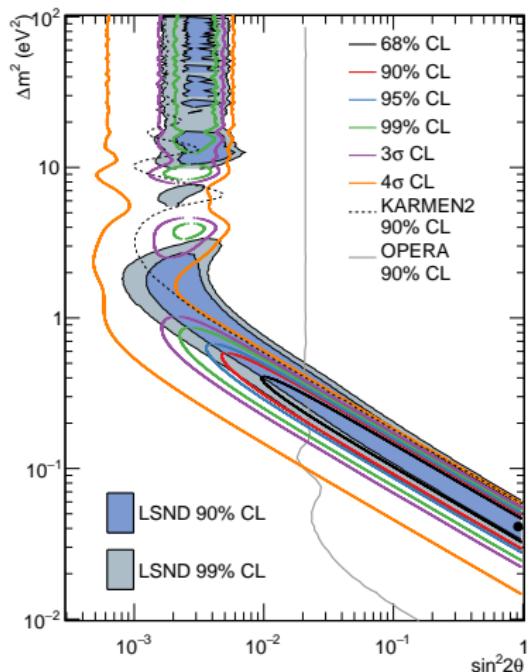
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[MiniBooNE, arxiv:1805.12028]



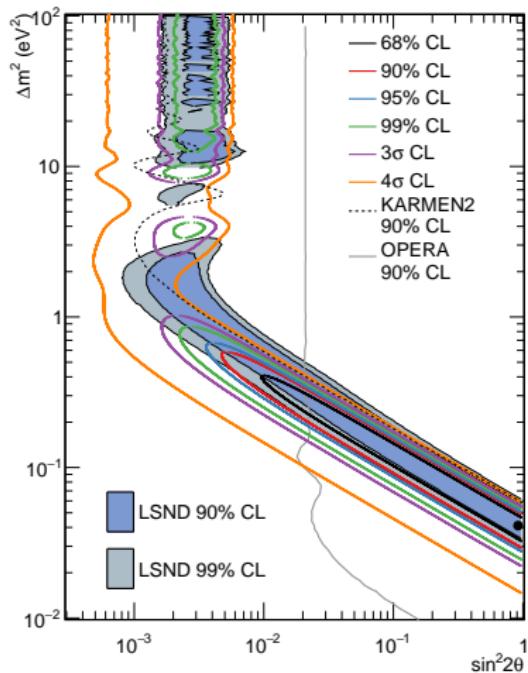
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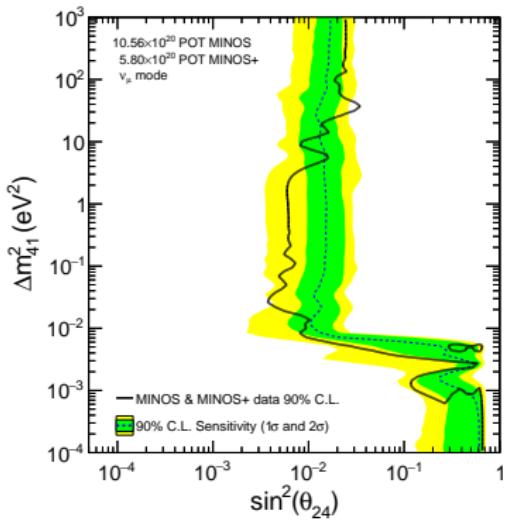


Recent results...

[MiniBooNE, arxiv:1805.12028]



[MINOS+, arxiv:1710.06488]



3+1 Neutrino Model

new $\Delta m_{\text{SBL}}^2 \Rightarrow 4$ neutrinos!



ν_4 with $m_4 \simeq 1$ eV,
no weak interactions



light sterile neutrino (LS ν)

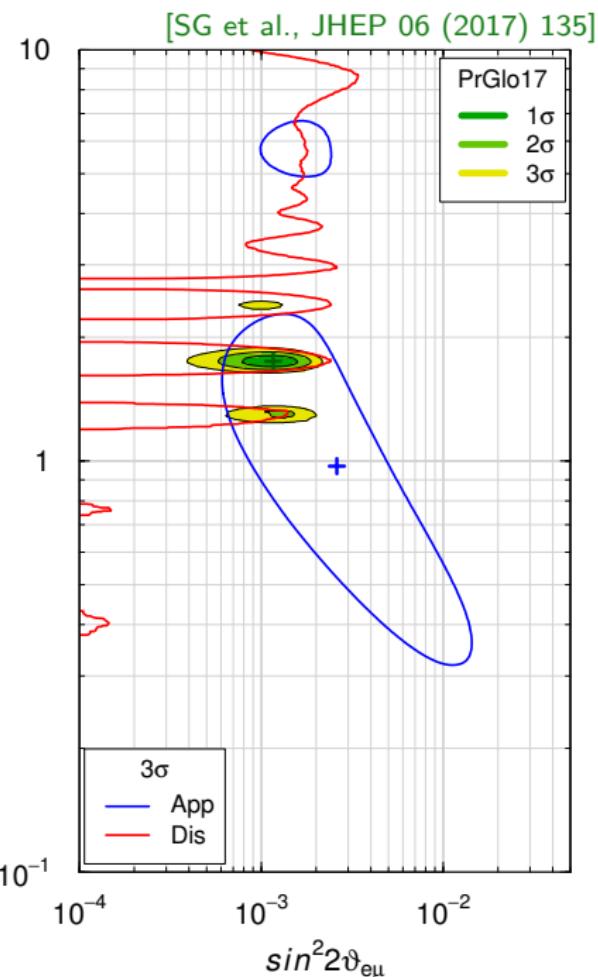
3 (active) + 1 (sterile) mixing:

$$\nu_\alpha = \sum_{k=1}^{3+1} U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau, s)$$

ν_s is mainly ν_4 :

$$m_s \simeq m_4 \simeq \sqrt{\Delta m_{41}^2} \simeq \sqrt{\Delta m_{\text{SBL}}^2}$$

assuming $m_4 \gg m_i$ ($i = 1, 2, 3$)



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[SG et al., in preparation (2018)]

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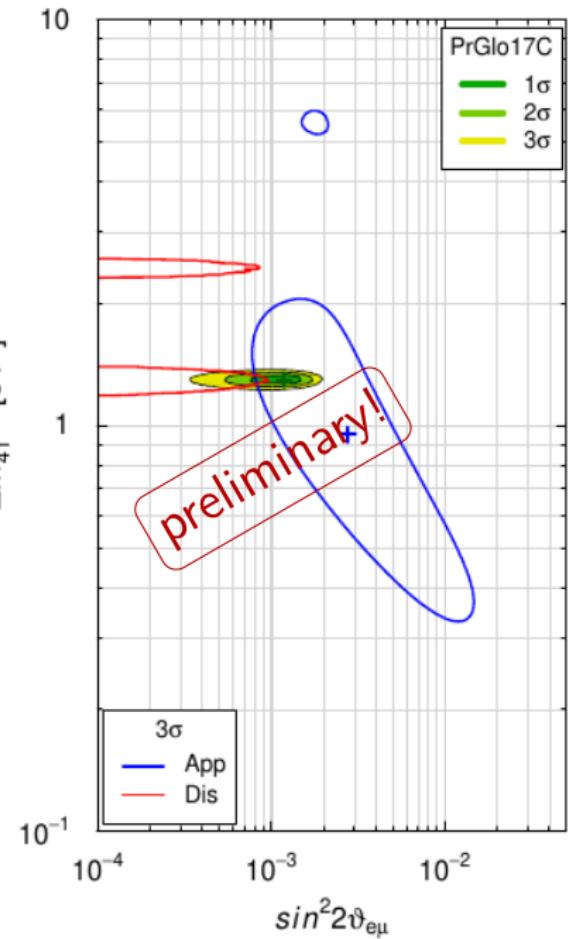
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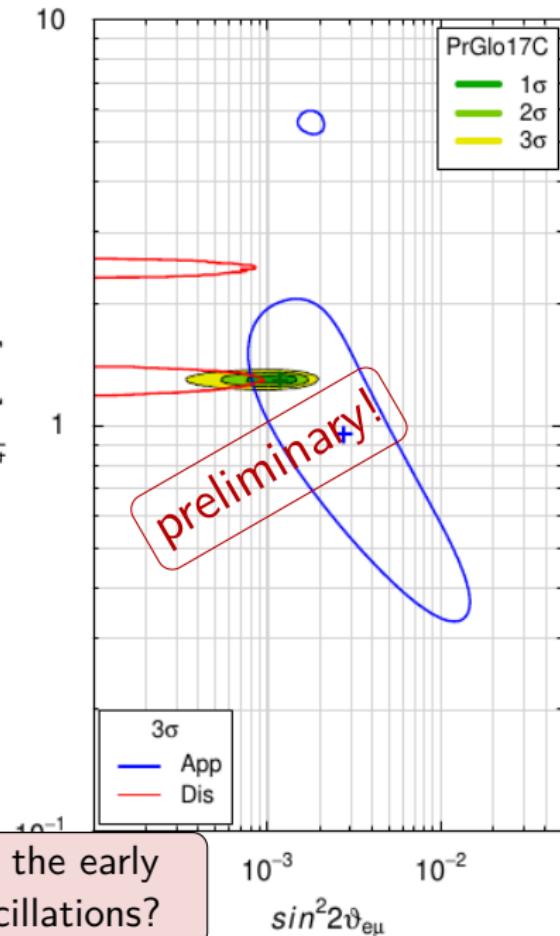
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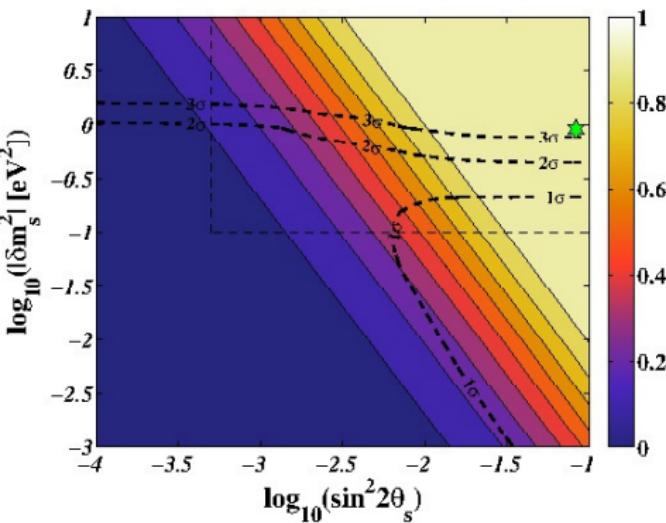


can ν_4 thermalize in the early
Universe through oscillations?

LS ν thermalization

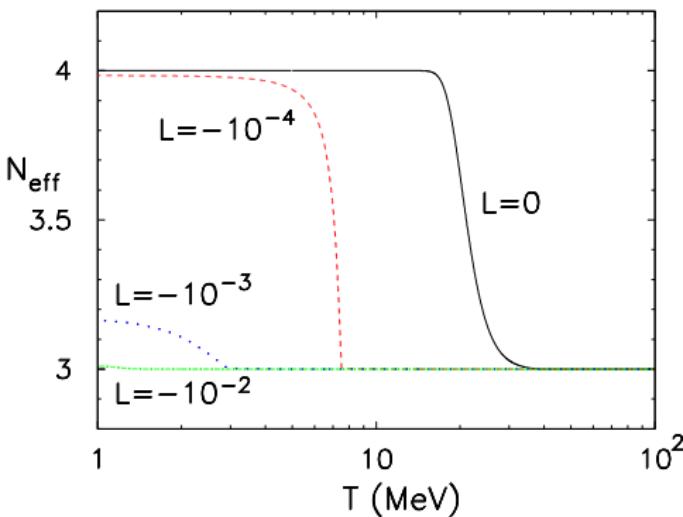
Using SBL best-fit parameters for the LS ν (Δm_{41}^2 , θ_s):

[Hannestad et al., JCAP 07 (2012) 025]



(colors coding ΔN_{eff})

[Mirizzi et al., PRD 86 (2012) 053009]



(L : lepton asymmetry)

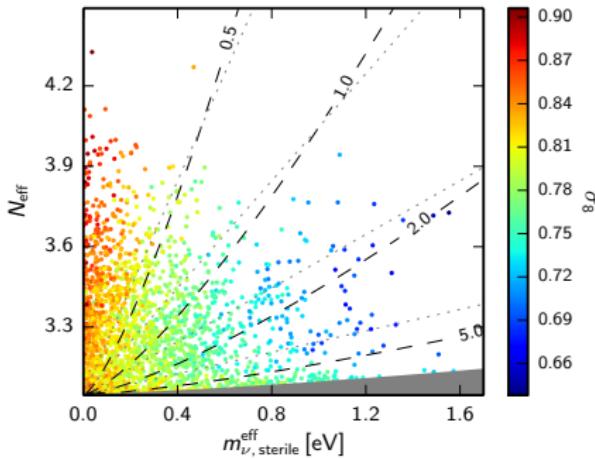
Unless $L \gtrsim \mathcal{O}(10^{-3})$, $\Delta N_{\text{eff}} \simeq 1$

See also: [Saviano et al., PRD 87 (2013) 073006], [Hannestad et al., JCAP 08 (2015) 019]

[to be precise: ΔN_{eff} is slightly smaller at CMB decoupling, when the LS ν starts to be non-relativistic]

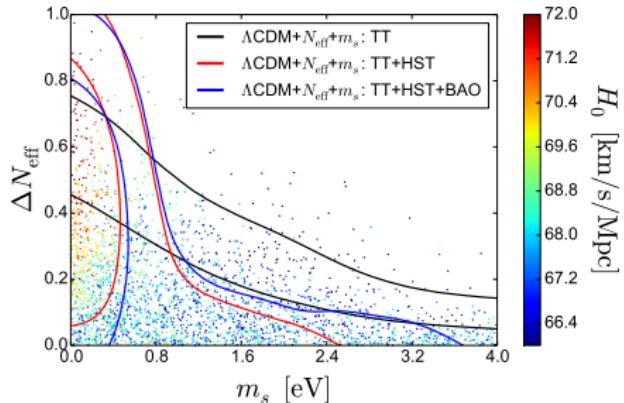
LS ν constraints from cosmology

CMB+local: [Planck Collaboration, 2015]



$$\left\{ \begin{array}{l} N_{\text{eff}} < 3.7 \\ m_s^{\text{eff}} < 0.52 \text{ eV} \end{array} \right. \quad \begin{array}{l} (\text{TT+lensing+BAO}) \\ [\text{TT+HST}] \end{array}$$

[Archidiacono et al., JCAP 08 (2016) 067]



dataset	free ΔN_{eff}	$\Delta N_{\text{eff}} = 1$
(TT)	$N_{\text{eff}} < 3.5$	$m_s < 0.66 \text{ eV}$
(+H ₀)	$N_{\text{eff}} < 3.9$	$m_s < 0.55 \text{ eV}$
(+BAO)	$N_{\text{eff}} < 3.8$	$m_s < 0.53 \text{ eV}$

BBN constraints: $N_{\text{eff}} = 2.90 \pm 0.22$ (BBN+Y_p) [Peimbert et al., 2016]

Summary: $\Delta N_{\text{eff}} = 1$ from LS ν incompatible with $m_s \simeq 1 \text{ eV}$!

Incomplete Thermalization

Active-sterile oscillations in the early Universe:

mixing parameters from SBL data $\Rightarrow \Delta N_{\text{eff}} \simeq 1$

[Hannestad et al., 2012] [Mirizzi et al., 2012]

Many probes constrain $\Delta N_{\text{eff}} < 1$. Do we need

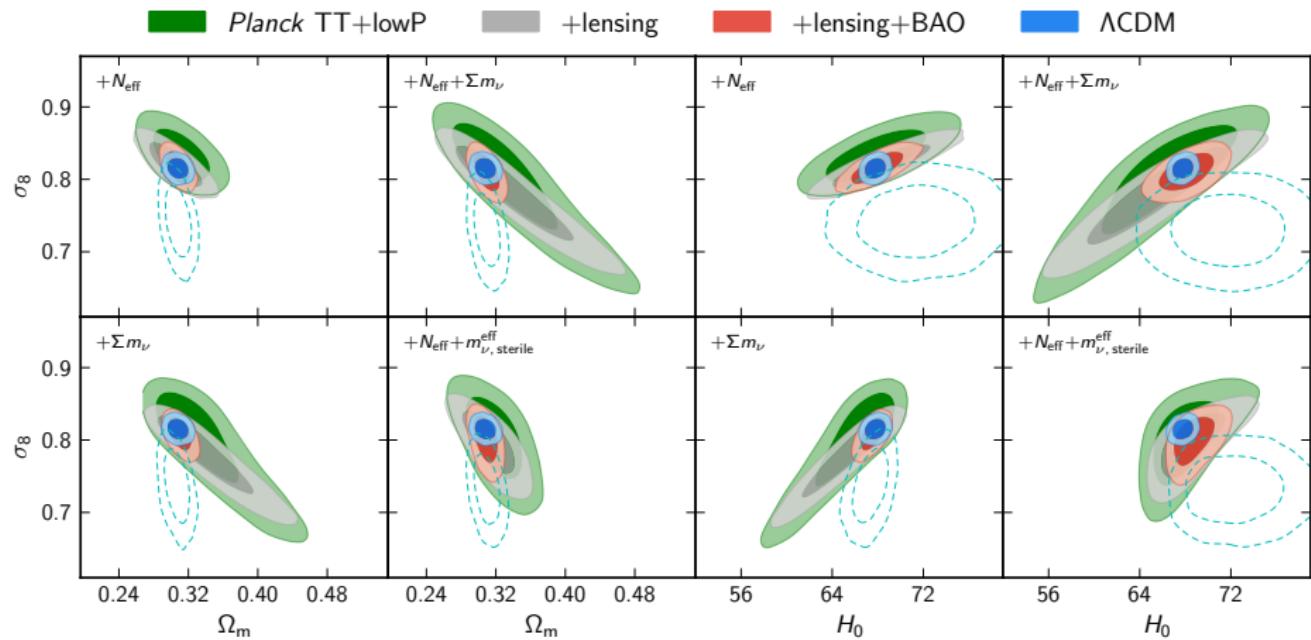
- a mechanism to suppress oscillations and full thermalization of ν_s ?
- to compensate $\Delta N_{\text{eff}} = 1$ with additional mechanisms in Cosmology?

Some ideas (an incomplete list!):

- large lepton asymmetry [Foot et al., 1995; Mirizzi et al., 2012; many more]
- new neutrino interactions [Bento et al., 2001; Dasgupta et al., 2014;
Hannestad et al., 2014; Saviano et al., 2014; Archidiacono et al. 2016; many more]
- entropy production after neutrino decoupling [Ho et al., 2013]
- very low reheating temperature [Gelmini et al., 2004; Smirnov et al., 2006]
- time varying dark energy components [Giusarma et al., 2012]
- larger expansion rate at the time of ν_s production [Rehagen et al., 2014]

Solving both σ_8 and H_0 Tension?

[Planck Collaboration, 2015]



dashed: local measurements – Λ CDM model, Λ CDM + $\nu_{a,s}$ models: full cosmological dataset

H_0 increases $\Rightarrow \sigma_8$ increases (and viceversa)!
The correlations do not help.

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The Strong CP problem

CP violation is permitted in QCD

effective periodic strong
CP-violating term Θ

Θ is an input of the SM

must be measured, cannot be predicted

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not observed! → Θ is small ($\lesssim 10^{-10}$)

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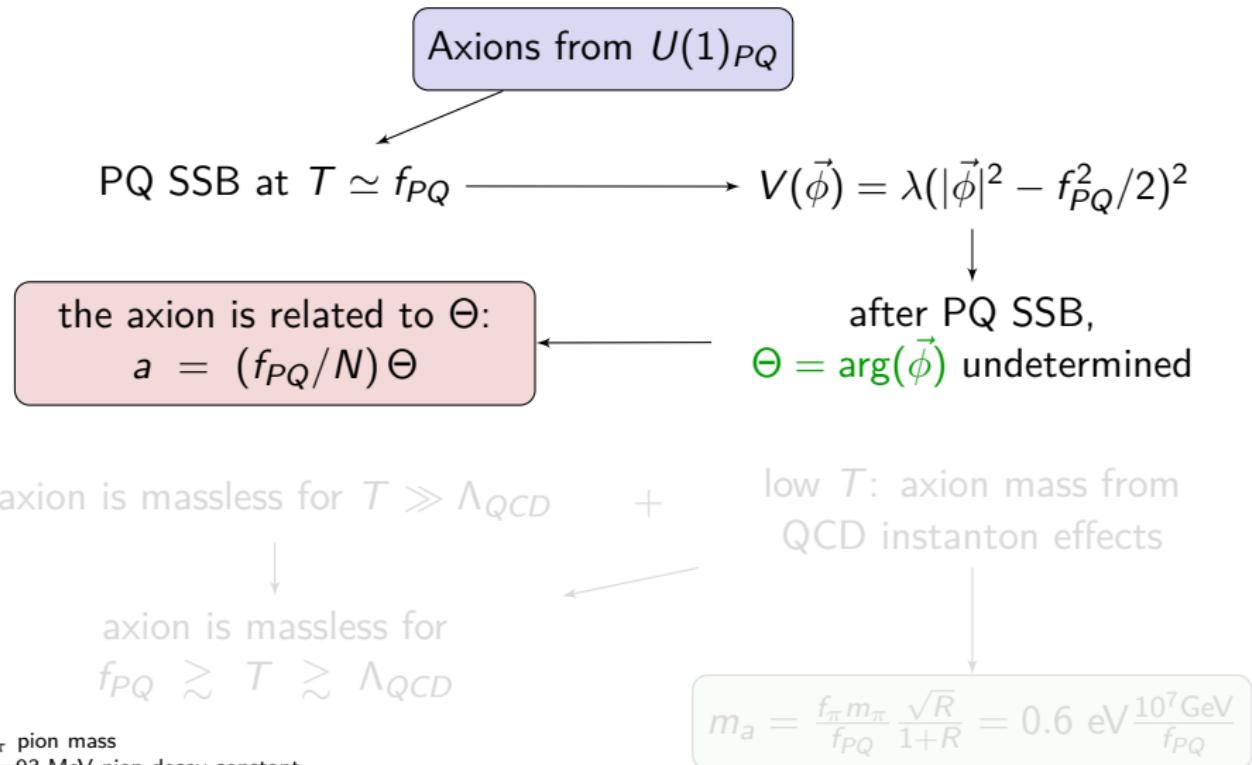
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Strong CP problem

[Peccei&Quinn,
1977]

Possible solution: Θ is a field, corresponding to
a new spontaneously broken symmetry $U(1)_{\text{PQ}}$

Axions during the expansion - I



m_π pion mass

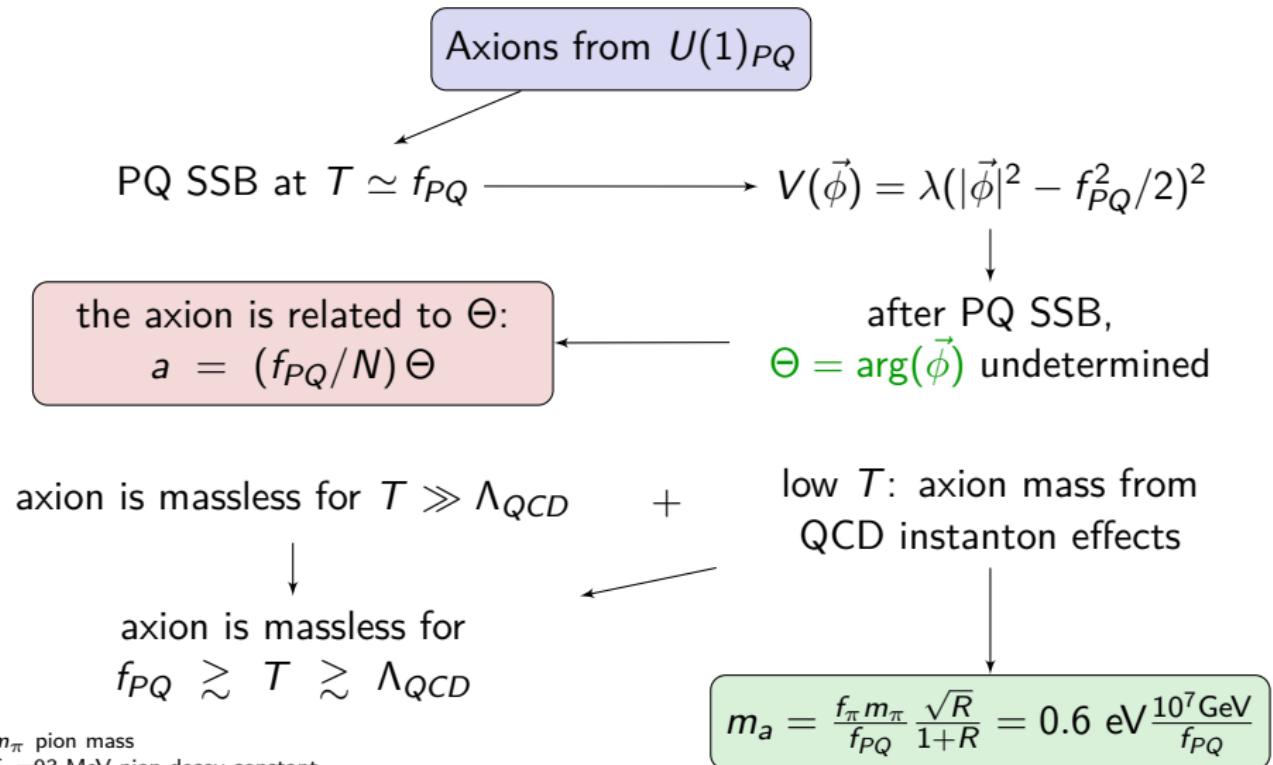
$f_\pi = 93$ MeV pion decay constant

N color anomaly of the PQ symmetry

$R = 0.553 \pm 0.043$ up-to-down quark masses ratio

PQ SSB = spontaneous symmetry breaking of PQ symmetry

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Axions during the expansion - II

[Archidiacono et al., JCAP 05 (2015) 050]

Note: axion couplings $\propto 1/f_{PQ} \propto m_a$

lighter axions interact less!

Axion production?

depends on interactions

thermal processes

non-thermal processes

misalignment

decay of axionic string

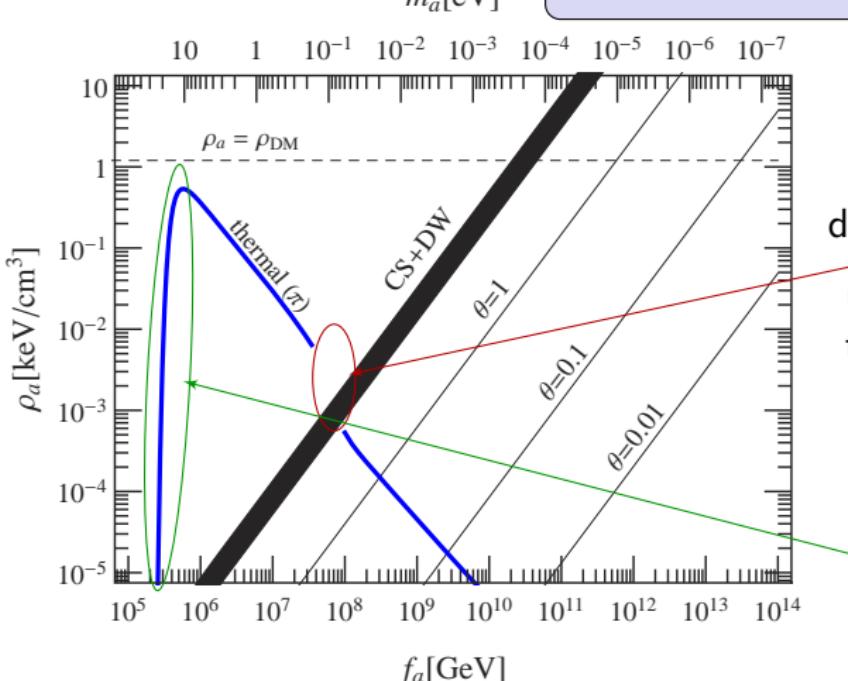
see also: [D. Marsh, Phys.Rept. 643 (2016) 1-79]

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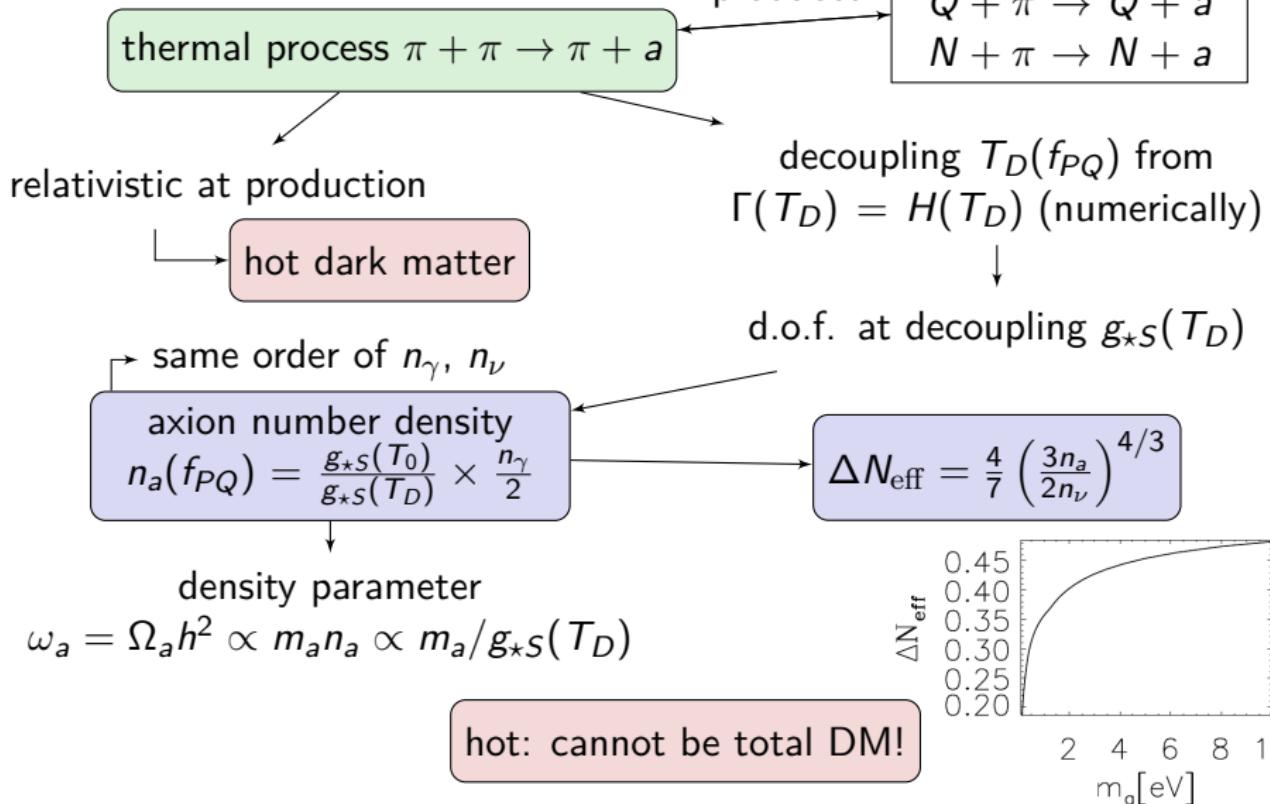
lighter axions interact less!



drop when $T_{fr} \simeq \Lambda_{QCD}$
(less efficient interactions with unconfined quarks and gluons + entropy dilution)

Axion rapidly decays $a \rightarrow 2\gamma$ and disappears

Thermal production



Q quarks, N nucleons

T_0 CMB temperature today

$n_\gamma (\nu)$ number density of photons (neutrinos)

S. Gariazzo

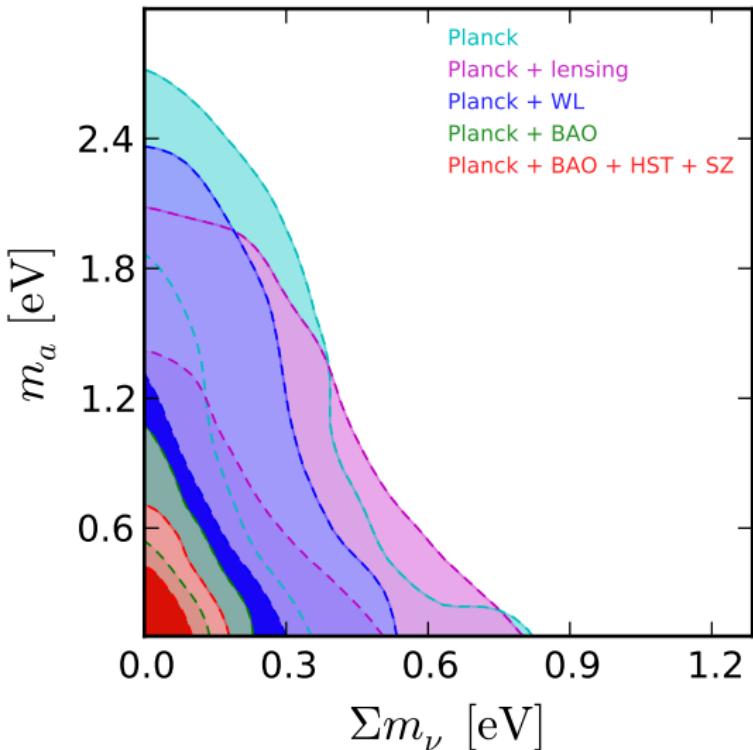
"Dark radiation candidates: light sterile neutrinos and thermal axions"

Ferrara, 27/06/2018

11/14

Constraints - I

[di Valentino et al., PLB 752 (2016) 182]



thermal axion behavior is similar to massive neutrinos



degeneracy
 $\sum m_\nu - m_a$

but different contributions to ΔN_{eff}

$(\Delta N_{\text{eff},a} \text{ depends on } m_a)$

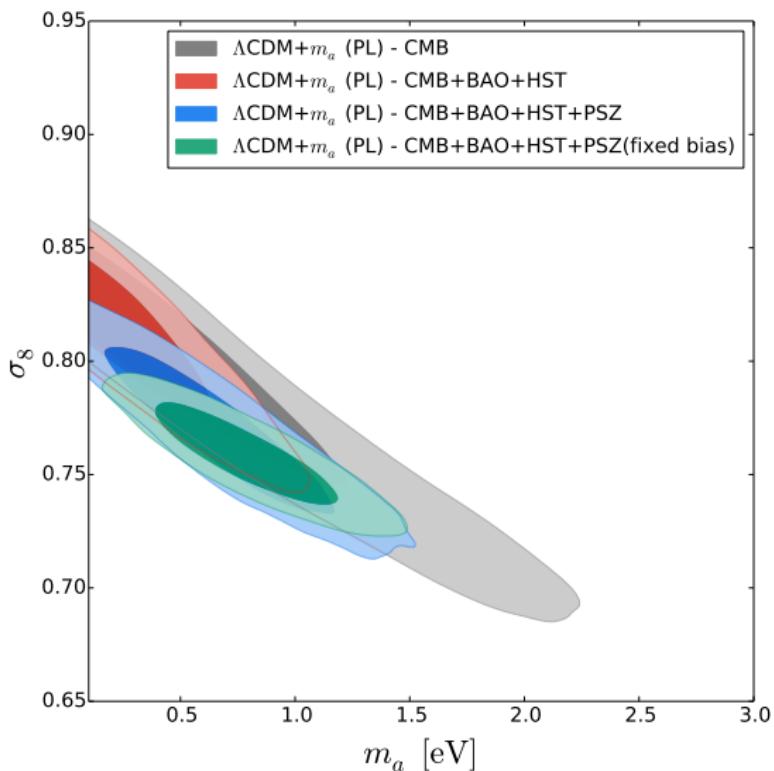


not complete degeneracy

Stronger constraint:
 $m_a < 0.529$ eV (95%)

Constraints - II

[Di Valentino et al., PRD 91 (2015) 123505]

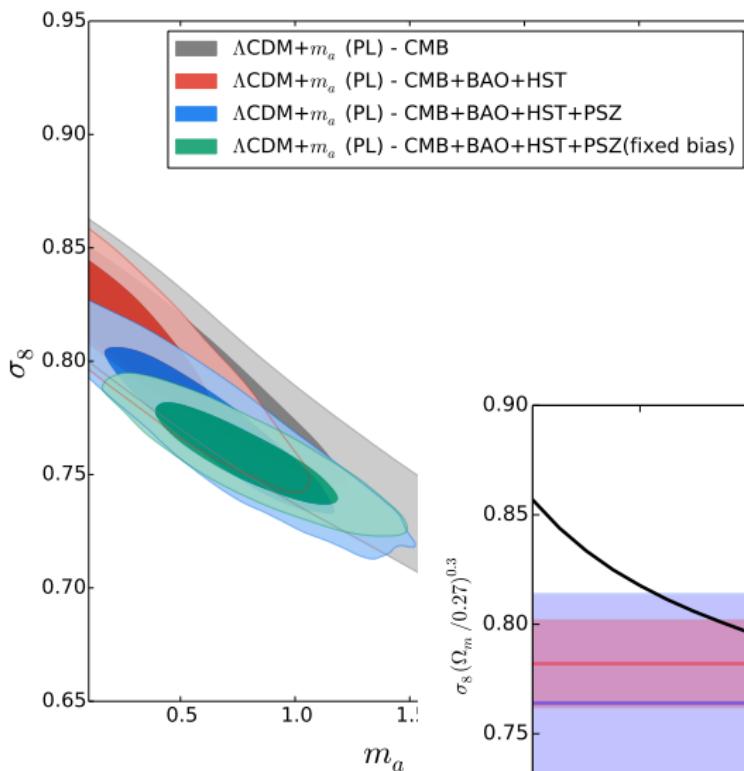


axion is hot relic
↓
suppression of matter power spectrum
↓
reduced fluctuations at small scales

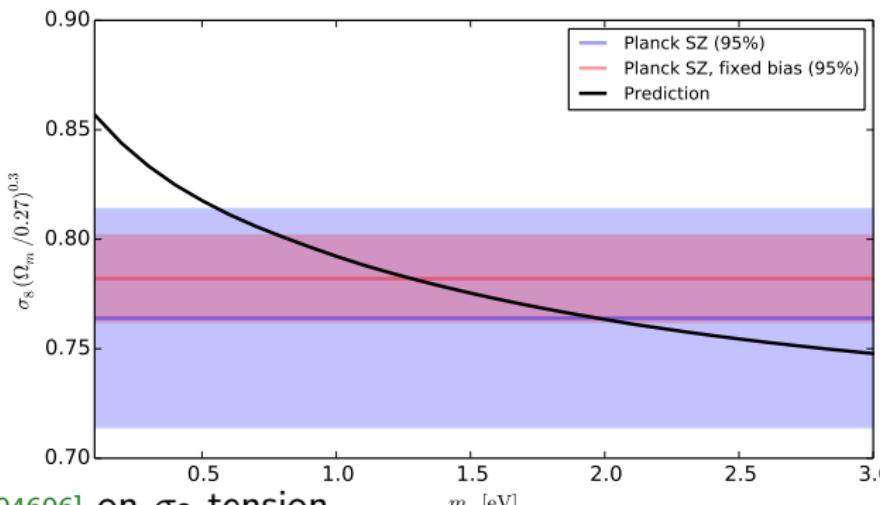
see also [Joudaki et al., 1610.04606] on σ_8 tension

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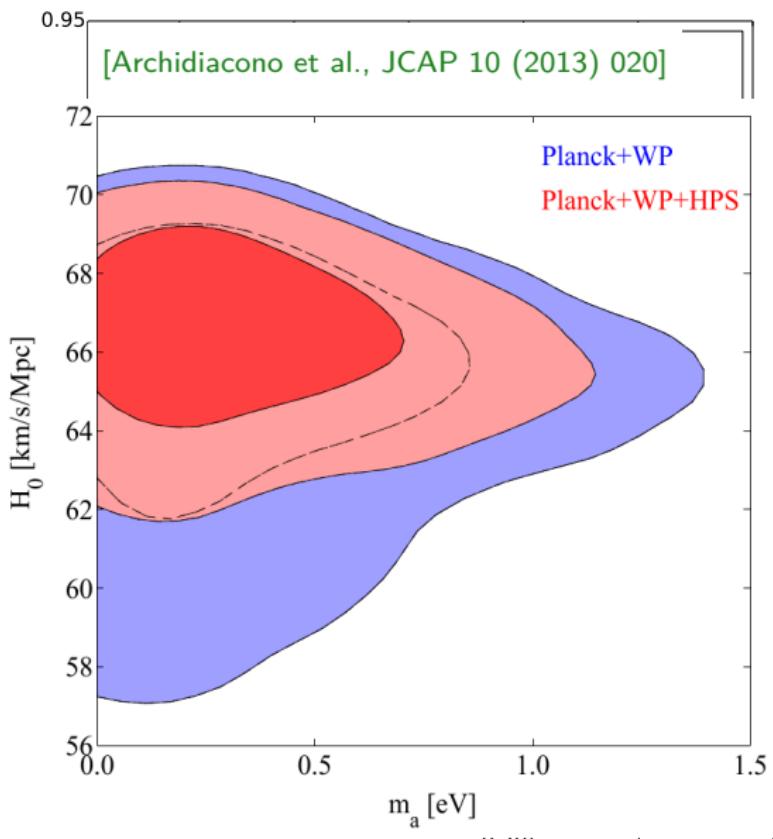
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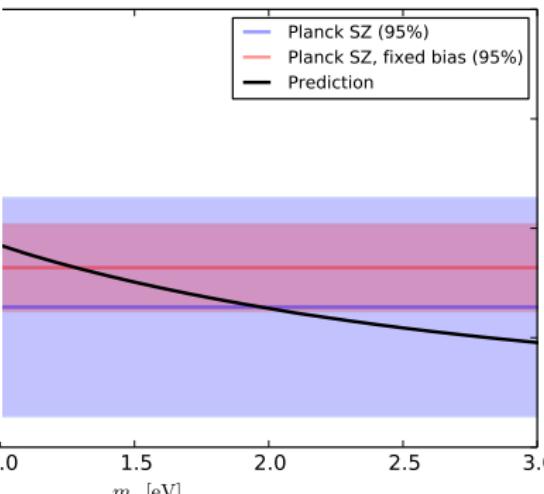
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Is there a light sterile neutrino?

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Not completely clear.

If yes, problems in early universe!

More new physics to be discovered?

What about the strong CP problem?

Is there an axion?

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If yes, study it using CMB data!

DM may be thermal + non-thermal axions/axion-like particles

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Dark radiation cannot really solve
the H_0 and σ_8 tensions. . .

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