#### The GSI Time Anomaly: Facts and Fiction

Carlo Giunti

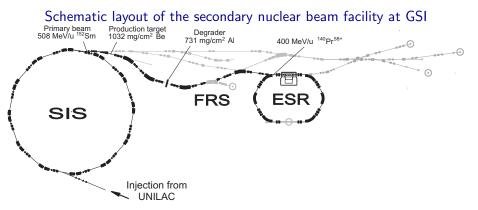
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#### **Outline**

- The GSI Anomaly
- Neutrino Mixing?
- Interference: Double-Slit Analogy
- Causality
- arXiv:0801.1465 and arXiv:0805.0435
- arXiv:0801.2121 arXiv:0801.3262
- Quantum Beats?
- Conclusions

## The GSI Experiment



{Litvinov et al, nucl-ex/0509019}

- SIS: Heavy Ion Synchrotron
- FRS: FRagment Separator
- ESR: Experiment Storage Ring

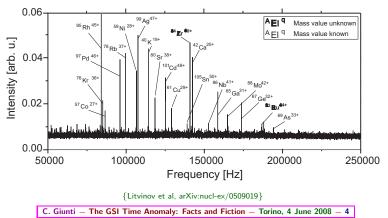


#### Schottky Mass Spectrometry

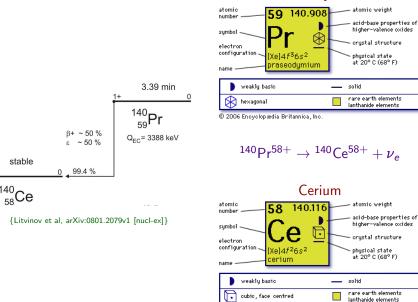
- $\blacktriangleright$  Stored ions circulate in ESR with revolution frequencies  $\sim 2 \mbox{ MHz}$
- At each turn they induce mirror charges on two electrodes
- Revolution frequency spectra provide information about q/m:

$$f = \frac{\omega}{2\pi} = \frac{qB}{2\pi m\gamma}$$

Area of each frequency peak is proportional to number of stored ions



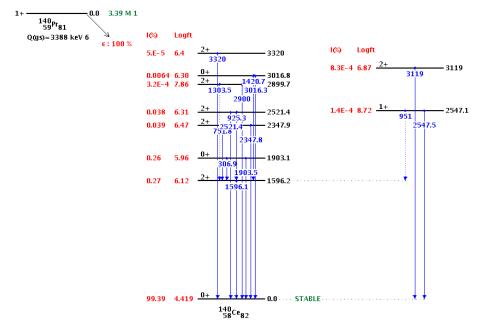
#### Praseodymium



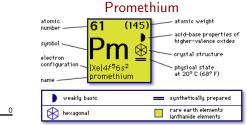
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140,

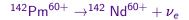
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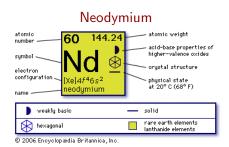


{NuDat, http://www.nndc.bnl.gov/nudat/}

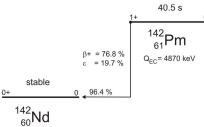


indicates the mass of the longest-lived isotope
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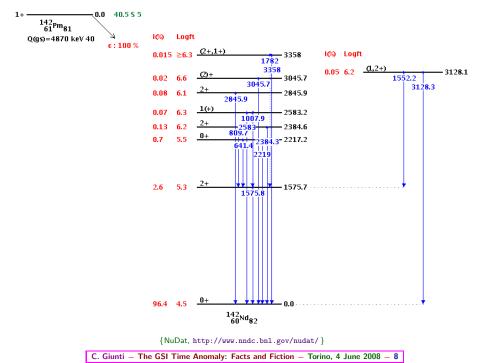


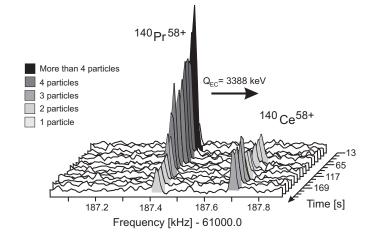




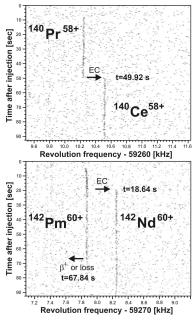


{Litvinov et al, arXiv:0801.2079v1 [nucl-ex]}





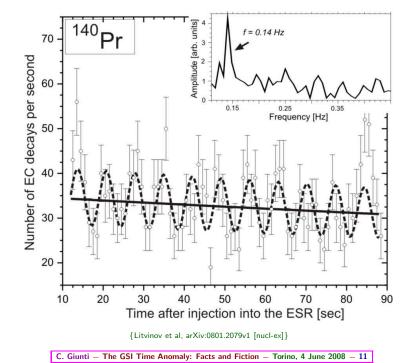
- About six initial <sup>140</sup> Pr<sup>58+</sup> ions (f =  $qB/2\pi m\gamma$ )
- ► Two decayed via nuclear electron capture into <sup>140</sup>Ce<sup>58+</sup>
- Seen because  $\Delta q = 0 \Rightarrow \Delta f/f = -\Delta m/m$  (small)
- Other decayed via β<sup>+</sup> decay (Δq = −1 ⇒ Δf ~ −150 kHz) or were lost (interactions with residual gas)

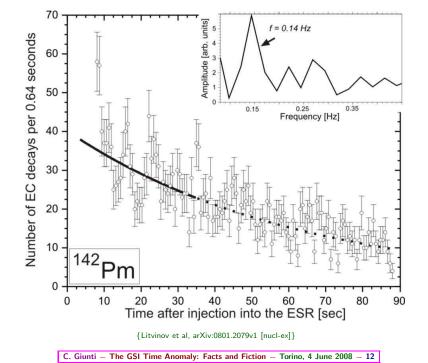


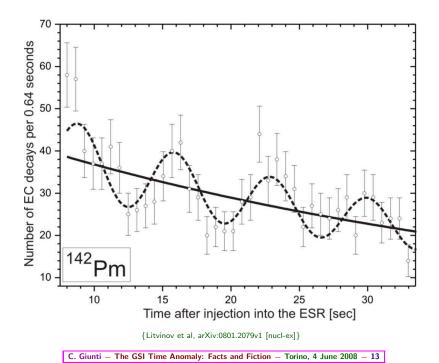
$$^{140}$$
Pr<sup>58+</sup>  $\rightarrow$   $^{140}$ Ce<sup>58+</sup> +  $\nu_e$ 

$$^{142}\mathsf{Pm}^{60+}
ightarrow^{142}\mathsf{Nd}^{60+}+
u_e$$

{Litvinov et al, arXiv:0801.2079v1 [nucl-ex]}







(1) 
$$\frac{dN_{EC}(t)}{dt} = \lambda_{EC} N(t) = \lambda_{EC} N(0) e^{-\lambda t}$$
  
(2) 
$$\frac{dN_{EC}(t)}{dt} = \tilde{\lambda}_{EC}(t) N(t) = \tilde{\lambda}_{EC}(t) N(0) e^{-\lambda t}$$

 $\lambda = \lambda_{\mathsf{EC}} + \lambda_{eta^+} + \lambda_{\mathsf{loss}}$   $\widetilde{\lambda}_{\mathsf{EC}}(t) = \lambda_{\mathsf{EC}} \left[ 1 + a \cos(\omega t + \phi) \right]$ 

Fit parameters of <sup>140</sup> Pr data					
Eq.	$N_0\lambda_{EC}$	λ	а	ω	$\chi^2/DoF$
(1)	34.9(18)	0.00138(10)	-	-	107.2/73
(2)	35.4(18)	0.00147(10)	0.18(3)	0.89(1)	67.18/70
Fit parameters of <sup>142</sup> Pm data					
Eq.	$N_0\lambda_{EC}$	λ	а	ω	$\chi^2/DoF$
(1)	46.8(40)	0.0240(42)	-	-	63.77/38
(2)	46.0(39)	0.0224(41)	0.23(4)	0.89(3)	31.82/35

 $T(^{140}Pr^{58+}) = 7.06 \pm 0.08 s$ 

 $T(^{142}Pm^{60+}) = 7.10 \pm 0.22 s$ 

 $\langle a \rangle = 0.20 \pm 0.02$ 

{Litvinov et al, arXiv:0801.2079v1 [nucl-ex]}

## **Neutrino Mixing?**

{Litvinov et al, arXiv:0801.2079v1 [nucl-ex]}

 $|I_i
ightarrow I_f+
u_e ~~|
u_e
angle=\cosartheta|
u_1
angle+\sinartheta|
u_2
angle$ 

Initial Ion: Momentum  $\vec{P} = 0$ , Energy E

Massive  $\nu_k$ : Momentum  $\vec{p}_k$ , Energy  $E_k = \sqrt{p_k^2 + m_k^2}$ 

Final Ion: Momentum  $-\vec{p}_k$ , Energy  $M + p_k^2/2M$ 

 $E_1 + M + p_1^2/2M = E$   $E_2 + M + p_2^2/2M = E$ 

$$\Delta E \equiv E_2 - E_1 \simeq \frac{\Delta m^2}{2M} \qquad \Delta m^2 \equiv m_2^2 - m_1^2$$

$$\Delta E \equiv E_2 - E_1 \simeq \frac{\Delta m^2}{2M}$$

$$\Delta m^2 = \Delta m_{\odot}^2 \simeq 8 \times 10^{-5} \text{ eV}^2 \qquad M \simeq 140 \text{ amu} \simeq 130 \text{ GeV}$$

$$\Delta E \simeq 3.1 \times 10^{-16} \text{ eV}$$

$$T = \frac{2\pi}{\Delta E} \gamma \simeq 19.1 \text{ s} \qquad \gamma = 1.43$$
about 3 times larger than  $T_{\text{GSI}} \simeq 7 \text{ s}$ 

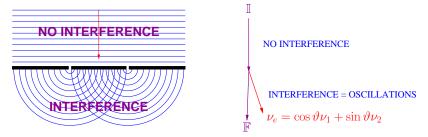
# Can the GSI Time Anomaly

## be due to

# Neutrino Mixing?

## NO

#### Interference: Double-Slit Analogy



- Decay rate of I corresponds to fraction of intensity of incoming wave which crosses the barrier
- Fraction of intensity of the incoming wave which crosses the barrier depends on the sizes of the holes
- It does not depend on interference effects which occur after the wave has passed through the barrier
- Analogy: decay rate of I cannot depend on interference of ν<sub>1</sub> and ν<sub>2</sub> which occurs after decay has happened



# INTERFERENCE OF $\nu_1$ AND $\nu_2$ OCCURRING AFTER THE DECAY CANNOT AFFECT THE DECAY RATE

#### arXiv:0801.1465 and arXiv:0805.0435

H.J. Lipkin

- Causality is violated explicitly
- ► arXiv:0801.1465: The difference in momentum  $\delta p_{\nu}$  between the two neutrino eigenstates with the same energy produces a small initial momentum change  $\delta P$  ...
- arXiv:0805.0435: Since the time dependence depends only on the propagation of the initial state, it is independent of the final state, which is created only at the decay point. Thus there is no violation of causality.
- ▶ But in calculation of effect: The phase difference at a time t between states produced by the neutrino mass difference on the motion of the initial ion in the laboratory frame with velocity V = (P/E) is

$$\delta\phi \approx -\delta E \cdot t = \Delta m^2/2E$$



#### <u>arXiv:0801.2121 – arXiv:0801.3262</u>

A. N. Ivanov, R. Reda, P. Kienle – M. Faber

•  $\mathbb{I} \to \mathbb{F} + \nu$  with final neutrino state  $|\nu\rangle = \sum_{k} |\nu_k\rangle$ 

Not even properly normalized to describe one particle:

$$\langle \nu_j | \nu_k 
angle = \delta_{jk} \implies \langle \nu | \nu 
angle = 3$$

Different from standard electron neutrino state

$$|
u_e
angle = \sum_k U_{ek}^* |
u_k
angle$$

► Standard QFT: 
$$P_{\mathbb{I} \to \mathbb{F} + \nu} = |\langle \nu, \mathbb{F} | \mathsf{S} | \mathbb{I} \rangle|^2 = \left| \sum_k \langle \nu_k, \mathbb{F} | \mathsf{S} | \mathbb{I} \rangle \right|^2$$

S-matrix operator at first order in perturbation theory:

$$\mathsf{S}=1-i\int\mathsf{d}^4x\,\mathscr{H}_W(x)$$

Effective four-fermion interaction Hamiltonian:

$$\mathcal{H}_{W}(x) = \frac{G_{F}}{\sqrt{2}} \cos \theta_{\mathsf{C}} \bar{\nu}_{e}(x) \gamma_{\rho} (1 - \gamma^{5}) e(x) \bar{n}(x) \gamma^{\rho} (1 - g_{\mathsf{A}} \gamma^{5}) p(x)$$
$$= \frac{G_{F}}{\sqrt{2}} \cos \theta_{\mathsf{C}} \sum_{k} U_{ek}^{*} \bar{\nu}_{k}(x) \gamma_{\rho} (1 - \gamma^{5}) e(x) \bar{n}(x) \gamma^{\rho} (1 - g_{\mathsf{A}} \gamma^{5}) p(x)$$

 $\langle \nu_{k}, \mathbb{F}|S|\mathbb{I} \rangle = U_{ek}^{*}\mathcal{M}_{k} \quad \text{with}$   $\mathcal{M}_{k} = -i\frac{G_{F}}{\sqrt{2}}\cos\theta_{C}\int d^{4}x \langle \nu_{k}, \mathbb{F}|\bar{\nu}_{k}(x)\gamma_{\rho}(1-\gamma^{5})e(x)\bar{n}(x)\gamma^{\rho}(1-g_{A}\gamma^{5})p(x)|\mathbb{I} \rangle$   $\mathsf{P}_{\mathbb{I} \to \mathbb{F}+\nu} = \left|\sum_{k} U_{ek}^{*}\mathcal{M}_{k}\right|^{2} \text{ different from standard } P = \sum_{k} |U_{ek}|^{2} |\mathcal{M}_{k}|^{2}$  [C. Giunti - The GSI Time Anomaly: Facts and Fiction - Torino, 4 June 2008 - 22]

Check: in the limit of massless neutrinos decay probability should reduce to the Standard Model decay probability

 $P_{\rm SM} = |\mathcal{M}_{\rm SM}|^2$ 

$$\mathcal{M}_{\mathsf{SM}} = -i \frac{G_F}{\sqrt{2}} \cos \theta_{\mathsf{C}} \int d^4 x \langle \nu_e, \mathbb{F} | \bar{\nu}_e(x) \gamma_\rho(1 - \gamma^5) e(x) \bar{n}(x) \gamma^\rho(1 - g_A \gamma^5) p(x) | \mathbb{I} \rangle$$

ما جارت

where  $\nu_e$  is the Standard Model massless electron neutrino  $\mathcal{M}_k = -i \frac{G_F}{\sqrt{2}} \cos\theta_C \int d^4 x \langle \nu_k, \mathbb{F} | \bar{\nu}_k(x) \gamma_\rho (1 - \gamma^5) e(x) \bar{n}(x) \gamma^\rho (1 - g_A \gamma^5) p(x) | \mathbb{I} \rangle$ 

$$\mathcal{M}_{k} \xrightarrow[m_{k} \to 0]{} \mathcal{M}_{\mathsf{SM}}$$
$$P_{\mathbb{I} \to \mathbb{F} + \nu} = \left| \sum_{k} U_{ek}^{*} \mathcal{M}_{k} \right|^{2} \xrightarrow[m_{k} \to 0]{} |\mathcal{M}_{\mathsf{SM}}|^{2} \left| \sum_{k} U_{ek}^{*} \right|^{2} \neq P_{\mathsf{SM}}$$

• Correct normalized final neutrino state ( $\langle \nu_e | \nu_e \rangle = 1$ ):

$$|\nu_{e}\rangle = \left(\sum_{j} |\langle \nu_{j}, \mathbb{F}|S|\mathbb{I}\rangle|^{2}\right)^{-1/2} \sum_{k} |\nu_{k}\rangle \langle \nu_{k}, \mathbb{F}|S|\mathbb{I}\rangle$$
$$= \left(\sum_{j} |U_{ej}|^{2} |\mathcal{M}_{j}|^{2}\right)^{-1/2} \sum_{k} U_{ek}^{*} \mathcal{M}_{k} |\nu_{k}\rangle$$

Standard decay probability:

$$P_{\mathbb{I} \to \mathbb{F} + \nu_{e}} = |\langle \nu_{e}, \mathbb{F} | \mathsf{S} | \mathbb{I} \rangle|^{2} = \sum_{k} |\langle \nu_{k}, \mathbb{F} | \mathsf{S} | \mathbb{I} \rangle|^{2} = \sum_{k} |U_{ek}|^{2} |\mathcal{M}_{k}|^{2}$$
$$P_{\mathbb{I} \to \mathbb{F} + \nu_{e}} \xrightarrow[m_{k} \to 0]{} P_{\mathsf{SM}}$$

 In experiments which are not sensitive to the differences of the neutrino masses, as neutrino oscillation experiments,

$$\mathcal{M}_k \simeq \overline{\mathcal{M}} \implies |\nu_e\rangle = \sum_k U_{ek}^* |\nu_k\rangle$$

#### **Quantum Beats?**

- ► GSI time anomaly can be due to interference effects in initial state
- ▶ Two coherent energy states of the decaying ion ⇒ Quantum Beats

 $\mathbb{I} = \mathcal{A}_1 \mathbb{I}_1 + \mathcal{A}_2 \mathbb{I}_2$ 

INTERFERENCE = QUANTUM BEATS

**INTERFERENCE = OSCILLATIONS**  $\nu_e = \cos \vartheta \nu_1 + \sin \vartheta \nu_2$ 

Incoming waves interfere at holes in barrier

INTERFERENCE

INTERFERENCE

Causality: interference due to different phases of incoming waves developed during propagation before reaching the barrier

- Quantum beats in GSI experiment can be due to interference of two coherent energy states of the decaying ion which develop different phases before the decay
- Coherence is preserved for a long time if measuring apparatus which monitors the ions with frequency ~ 2 MHz does not distinguish between the two states

$$\begin{split} |\mathbb{I}(t=0)\rangle &= \mathcal{A}_{1} |\mathbb{I}_{1}\rangle + \mathcal{A}_{2} |\mathbb{I}_{2}\rangle \qquad (|\mathcal{A}_{1}|^{2} + |\mathcal{A}_{2}|^{2} = 1) \\ \Gamma &= \Gamma_{1} \simeq \Gamma_{2} \implies |\mathbb{I}(t)\rangle = \left(\mathcal{A}_{1} e^{-iE_{1}t} |\mathbb{I}_{1}\rangle + \mathcal{A}_{2} e^{-iE_{2}t} |\mathbb{I}_{2}\rangle\right) e^{-\Gamma t/2} \\ P_{\mathsf{EC}}(t) &= |\langle \nu_{e}, \mathbb{F}|\mathsf{S}|\mathbb{I}(t)\rangle|^{2} = [1 + A\cos(\Delta Et + \varphi)] \overline{P}_{\mathsf{EC}} e^{-\Gamma t} \\ \mathcal{A} &\equiv 2|\mathcal{A}_{1}||\mathcal{A}_{2}|, \quad \Delta E \equiv E_{2} - E_{1}, \quad \overline{P}_{\mathsf{EC}} = |\langle \nu_{e}, \mathbb{F}|\mathsf{S}|\mathbb{I}_{1}\rangle|^{2} \simeq |\langle \nu_{e}, \mathbb{F}|\mathsf{S}|\mathbb{I}_{2}\rangle|^{2} \\ &= \frac{dN_{\mathsf{EC}}(t)}{dt} = N(0) \left[1 + A\cos(\Delta Et + \varphi)\right] \overline{\Gamma}_{\mathsf{EC}} e^{-\Gamma t} \end{split}$$

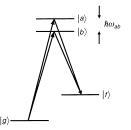


Fig. 1 Diagram of the four level system. A photon is absorbed by the ground state  $|b\rangle$  and excites a superposition of states  $|a\rangle$  and  $|b\rangle$  whose energy separation is  $\Delta E = \hbar \omega_{ab}$ . Emission of a second photon leaves the system in the final state  $|f\rangle$ .

$$\begin{split} I_{\rm fl}(t) &\propto \left| \mu_{ag} \right|^2 \left| \mu_{fa} \right|^2 {\rm e}^{-\gamma_{a}t} + \left| \mu_{bg} \right|^2 \left| \mu_{fb} \right|^2 {\rm e}^{-\gamma_{a}t} \\ &+ \left| \mu_{ag} \mu_{bg} \mu_{fa} \mu_{fb} \right| {\rm e}^{-(\gamma_{a}+\gamma_{b})t/2} \cos(\omega_{ab}t + \theta). \end{split}$$
(4)

Examination of this expression shows that it consists of two parts, one incoherent term (first two terms) describing the independent decays of the two states  $|a\rangle$  and  $|b\rangle$  and one coherent or cross term (last term) which decays at the average rate of the two states and, most importantly, is modulated at the angular frequency  $\omega_{ab}$ . The modulation frequency is the difference of the two angular frequencies in eqn. (2), *i.e.*  $\omega_{ab} =$  $|\omega_a - \omega_b|$ , and the coherent term in eqn. (4) is therefore termed the quantum beat. The angle  $\theta$  is included in eqn. (4) to describe the phase of the quantum beat, which depends on a number of factors such as the excitation and detection polarisations and transitions. When the transition moments and decay rates are equal, as is often the case, a particularly simple expression is derived for the four level system. In this case eqn. (4) becomes

$$I_{fl}(t) \propto [1 + \cos(\omega_{ab}t + \theta)]e^{-\gamma t}$$
, (5)

clearly illustrating the contributions of the incoherent and coherent terms to the fluorescence decay. In this special case the quantum beat is 100% modulated. It is important to point out

{Carter, Huber, Quantum beat spectroscopy in chemistry, Chem. Soc. Rev., 29 (2000) 305}

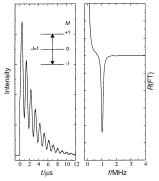


Fig. 2Zeeman quantum beat recorded for the R(0) line of the 17U transition in CS<sub>2</sub> in an external field of ~15 Gauss. The laser polarisation was perpendicular to the magnetic field direction and prepares a coherence between the  $M = \pm 1$  sublevels as shown in the level diagram. This is manifested by a single quantum beat on the fluorescence decay: the real part of the Fourier transform is also shown. The less than 100% modulation, which is observed in virtually all quantum beat measurements in molecules, is due to incoherent emission from the excited states.

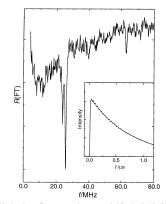


Fig. 6 Nuclear hyperfine quantum beats recorded for the  $P_{21}(Q_1(32))$  line in a vibrational band of the  $A^{+2}r - X^{-1}I$  transition in the  $A^{-1}OD$  van der Waals complex. The inset shows the fluorescence decay which exhibits weakly modulated quantum beats. Following Fourier transformation the beat frequencies between hyperfine levels in the  $A^{+2}r$  state are clearly visible.

{Carter, Huber, Quantum beat spectroscopy in chemistry, Chem. Soc. Rev., 29 (2000) 305}

$$\frac{\mathrm{d}N_{\mathsf{EC}}(t)}{\mathrm{d}t} = N(0) \left[1 + A\cos(\Delta E t + \varphi)\right] \overline{\mathsf{\Gamma}}_{\mathsf{EC}} e^{-\mathsf{\Gamma} t}$$

 $\Delta E(^{140} \mathrm{Pr}^{58+}) = (5.86 \pm 0.07) \times 10^{-16} \,\mathrm{eV}\,, \qquad A(^{140} \mathrm{Pr}^{58+}) = 0.18 \pm 0.03$ 

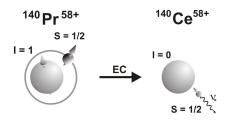
 $\Delta E(^{142} Pm^{60+}) = (5.82 \pm 0.18) \times 10^{-16} \text{ eV}, \quad A(^{142} Pm^{60+}) = 0.23 \pm 0.04$ 

 $A \equiv 2|\mathcal{A}_1||\mathcal{A}_2|$ 

- Energy splitting is extremely small
- $|\mathcal{A}_1|^2/|\mathcal{A}_2|^2 \sim 1/99$  or  $|\mathcal{A}_2|^2/|\mathcal{A}_1|^2 \sim 1/99$
- It is difficult to find an appropriate mechanism

## Hyperfine Splitting

smallest known energy splitting



 $F = I + s - \begin{pmatrix} 3/2 & \longrightarrow \\ 1/2 & \longrightarrow \end{pmatrix} F = I + s = 1/2$ 

{Litvinov et al, PRL 99 (2007) 262501, arXiv:0711.3709}

 $\begin{array}{rcl} \Delta E \sim 10^{-6} \, \mathrm{eV} & \Longrightarrow & T \sim 10^{-9} \, \mathrm{s} & \mathrm{f} \sim \mathrm{GHz} \\ & & & & & & & & \\ & & & & & & & \\ T_{\mathrm{GSI}} \simeq 7 \, \mathrm{s} & & & & & \\ f_{\mathrm{GSI}} \simeq 0.14 \mathrm{Hz} & & & & & & \\ \Delta E_{\mathrm{GSI}} = 2\pi/T_{\mathrm{GSI}} \simeq 6 \times 10^{-16} \, \mathrm{eV} \end{array}$ 

#### **Conclusions**

- Interference: due to phase difference of two incoming waves
- Causality: there cannot be interference of waves before they exist
- The GSI ion lifetime anomaly cannot be due to interference of decay product before the decay product start to exist (neutrino mixing in the final state)
- The GSI ion lifetime anomaly can be due to interference of two energy states of the decaying ion: Quantum Beats
- ► No known mechanism, because
  - Energy splitting of the two energy states:  $\Delta E \sim 6 \times 10^{-16} \, \mathrm{eV}$
  - Ratio of probabilities of the two energy states: 1/99