Sterile Neutrinos: Phenomenology and Fits

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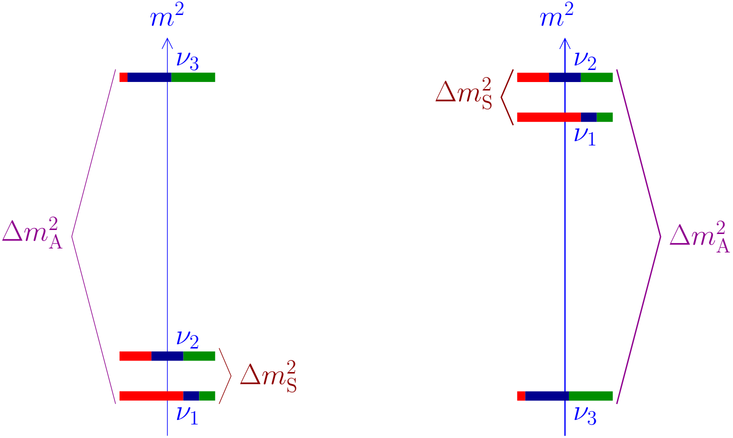
Three-Neutrino Mixing Paradigm

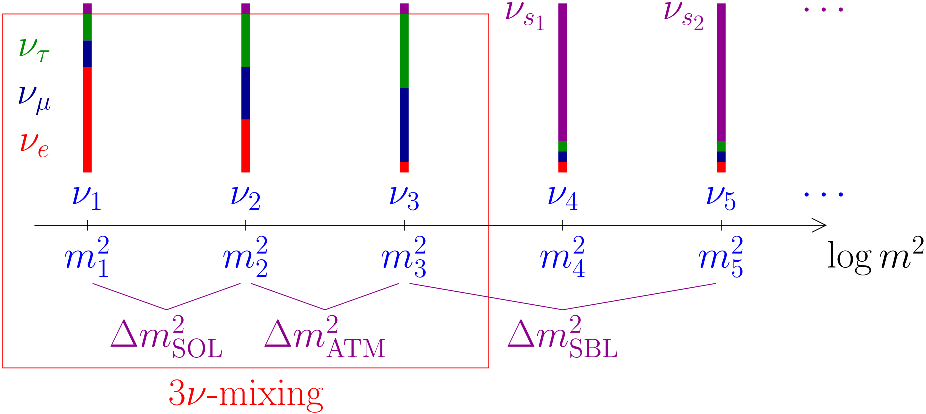
ν_e ν_μ ν_τ



Normal Spectrum

Inverted Spectrum

Beyond Three-Neutrino Mixing



Standard Model

- Neutrinos are the only massless fermions
- Neutrinos are the only fermions with only left-handed component ν_L
- Neutrinos are the only neutral fermions

Extension of the SM: Massive Neutrinos

- Simplest extension: introduce right-handed component ν_R (singlet of $SU(2)_L \times U(1)_Y$)
- One generation: Dirac mass $m_{\text{D}}\overline{\nu_R}\nu_L$ + Majorana mass $m_{\text{M}}\overline{\nu_R^c}\nu_R$
  $\Longrightarrow$ 2 massive Majorana neutrinos
- Three left-handed fields + N_R right-handed fields:
  $\nu_{eL}\,,\ \nu_{\mu L}\,,\ \nu_{\tau L}\ +\ \nu_{1R}\,,\ \ldots\,,\ \nu_{N_R R}$
  $\Longrightarrow$ $3 + N_R$ massive Majorana neutrinos

Sterile Neutrinos

- Light anti-ν_R are called sterile neutrinos

$$\nu_R^c \to \nu_{sL} \qquad \text{(left-handed)}$$

- Sterile means no standard model interactions
- Active neutrinos $(\nu_e, \nu_\mu, \nu_\tau)$ can oscillate into sterile neutrinos (ν_s)
- Observables:
  - Disappearance of active neutrinos (neutral current deficit)
  - Indirect evidence through combined fit of data (current indication)
- Short-baseline anomalies + 3ν-mixing:

$$\begin{array}{ccccc} & \Delta m^2_{21} \ll & |\Delta m^2_{31}| \ll & |\Delta m^2_{41}| \leq & \ldots \\ \nu_1 & \nu_2 & \nu_3 & \nu_4 & \ldots \\ \nu_e & \nu_\mu & \nu_\tau & \nu_{s_1} & \ldots \end{array}$$

- In this talk I consider sterile neutrinos with mass scale $\sim 1\,\mathrm{eV}$ in light of LSND, MiniBooNE, Reactor Anomaly, Gallium Anomaly.
- Other possibilities (not incompatible):
  - Very light sterile neutrinos with mass scale $\ll 1\,\mathrm{eV}$: important for solar neutrino phenomenology

    [de Holanda, Smirnov, PRD 83 (2011) 113011]
  - Heavy sterile neutrinos with mass scale $\gg 1\,\mathrm{eV}$: could be Warm Dark Matter

    [Kusenko, Phys. Rept. 481 (2009) 1]

    [Boyarsky, Ruchayskiy, Shaposhnikov, Ann. Rev. Nucl. Part. Sci. 59 (2009) 191]

Effective SBL Oscillation Probabilities in 3+1 Schemes

$$P_{\nu_\alpha \to \nu_\beta} = \sin^2 2\vartheta_{\alpha\beta} \sin^2 \left(\frac{\Delta m^2_{41} L}{4E} \right) \qquad \sin^2 2\vartheta_{\alpha\beta} = 4 |U_{\alpha 4}|^2 |U_{\beta 4}|^2$$

No CP Violation!

$$P_{\nu_\alpha \to \nu_\alpha} = 1 - \sin^2 2\vartheta_{\alpha\alpha} \sin^2 \left(\frac{\Delta m^2_{41} L}{4E} \right) \qquad \sin^2 2\vartheta_{\alpha\alpha} = 4 |U_{\alpha 4}|^2 \left(1 - |U_{\alpha 4}|^2 \right)$$

Perturbation of 3ν Mixing

$$|U_{e4}|^2 \ll 1 \,, \quad |U_{\mu 4}|^2 \ll 1 \,, \quad |U_{\tau 4}|^2 \ll 1 \,, \quad |U_{s4}|^2 \simeq 1$$

$$U = \left(\begin{array}{ccc|c|} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{array} \right)$$

↑

SBL

$$\sin^2 2\vartheta_{\alpha\alpha} \ll 1$$

$$\Downarrow$$

$$|U_{\alpha 4}|^2 \simeq \frac{\sin^2 2\vartheta_{\alpha\alpha}}{4}$$

Effective SBL Oscillation Probabilities in 3+2 Schemes

$$\phi_{kj} = \Delta m^2_{kj} L/4E$$

$$\eta = \arg[U^*_{e4} U_{\mu4} U_{e5} U^*_{\mu5}]$$

$$\begin{aligned} P_{\overset{(-)}{\nu}_\mu \to \overset{(-)}{\nu}_e} = {} & 4|U_{e4}|^2|U_{\mu4}|^2 \sin^2\phi_{41} + 4|U_{e5}|^2|U_{\mu5}|^2 \sin^2\phi_{51} \\ & + 8|U_{\mu4}U_{e4}U_{\mu5}U_{e5}| \sin\phi_{41} \sin\phi_{51} \cos(\phi_{54} \overset{(+)}{-} \eta) \end{aligned}$$

$$\begin{aligned} P_{\overset{(-)}{\nu}_\alpha \to \overset{(-)}{\nu}_\alpha} = {} & 1 - 4(1 - |U_{\alpha4}|^2 - |U_{\alpha5}|^2)(|U_{\alpha4}|^2 \sin^2\phi_{41} + |U_{\alpha5}|^2 \sin^2\phi_{51}) \\ & - 4|U_{\alpha4}|^2|U_{\alpha5}|^2 \sin^2\phi_{54} \end{aligned}$$

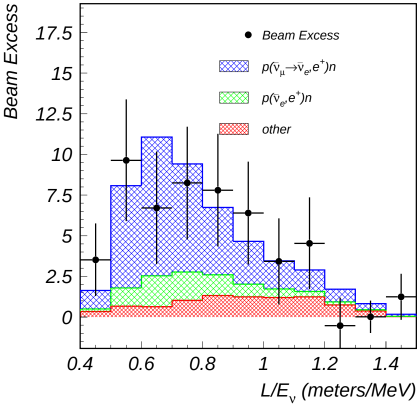
- More parameters: 7 (vs 3 in 3+1)
- CP violation

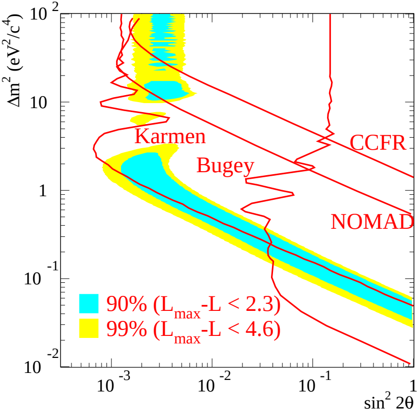
[Sorel, Conrad, Shaevitz, PRD 70 (2004) 073004; Maltoni, Schwetz, PRD 76 (2007) 093005; Karagiorgi et al, PRD 80 (2009) 073001; Kopp, Maltoni, Schwetz, PRL 107 (2011) 091801; Giunti, Laveder, PRD 84 (2011) 073008; Donini et al, arXiv:1205.5230]

LSND

[LSND, PRL 75 (1995) 2650; PRC 54 (1996) 2685; PRL 77 (1996) 3082; PRD 64 (2001) 112007]

$\bar{\nu}_\mu \to \bar{\nu}_e$ $\qquad L \simeq 30\,\text{m}$ $\qquad 20\,\text{MeV} \leq E \leq 200\,\text{MeV}$



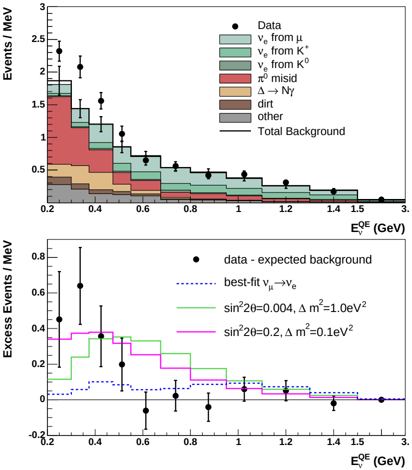


$$\Delta m^2_{\text{LSND}} \gtrsim 0.2\,\text{eV}^2 \qquad (\gg \Delta m^2_{\text{A}} \gg \Delta m^2_{\text{S}})$$

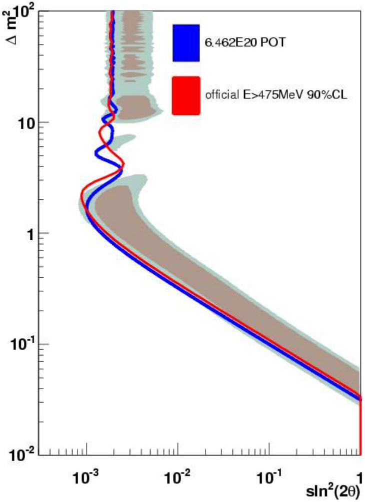
MiniBooNE Neutrinos

[PRL 98 (2007) 231801; PRL 102 (2009) 101802]

$\nu_\mu \to \nu_e$ $\quad L \simeq 541\,\text{m}$ $\quad 200\,\text{MeV} \le E \lesssim 3\,\text{GeV}$



[MiniBooNE, PRL 102 (2009) 101802]



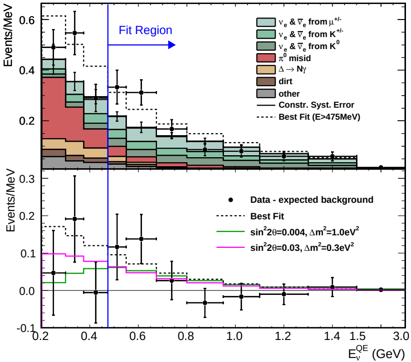
[Djurcic, arXiv:0901.1648]

- no $\nu_\mu \to \nu_e$ signal corresponding to LSND $\bar{\nu}_\mu \to \bar{\nu}_e$ signal ($E > 475\,\text{MeV}$)
- low-energy anomaly

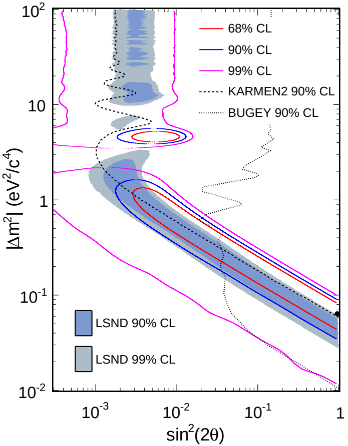
MiniBooNE Antineutrinos - 2009-2010

[PRL 103 (2009) 111801; PRL 105 (2010) 181801]

$\bar{\nu}_\mu \to \bar{\nu}_e$ $L \simeq 541\,\mathrm{m}$ $200\,\mathrm{MeV} \leq E \lesssim 3\,\mathrm{GeV}$

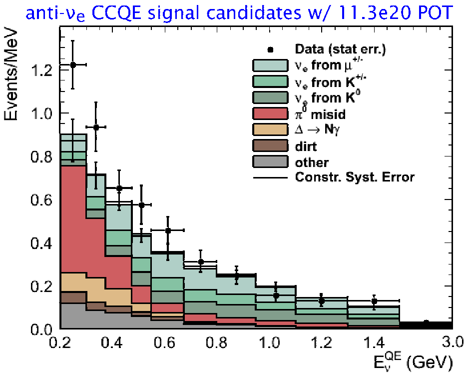


[MiniBooNE, PRL 105 (2010) 181801]



- 5.7e20 POT: agreement with LSND $\bar{\nu}_\mu \to \bar{\nu}_e$ signal ($E > 475\,\mathrm{MeV}$)
- similar L/E but different L and E $\Longrightarrow$ oscillations

MiniBooNE $\bar{\nu}$ - Neutrino 2012 - 6 June



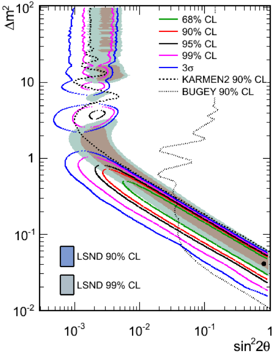
Higher stat anti-neutrino data is now much more consistent with what was observed in the data taken with a neutrino beam

* Systematic error after all other data constraints applied, e.g. ν_μ CCQE, NC π^0 , dirt events, SciBooNE K^+

		1st half			2nd half	
	data	mc	excess	data	mc	excess
200-475	119	100.5±14.3	18.5 (1.3s)	138	100.0±14.1	38 (2.7s)
475-1250	120	99.1±14.0	20.9 (1.5s)	101	103.1±14.4	-2.2 (-0.2s)

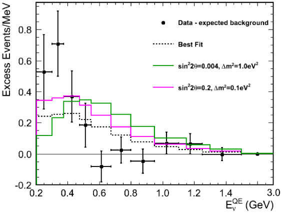
agreement with LSND signal is sadly vanishing

MiniBooNE ν and $\bar{\nu}$ - Neutrino 2012 - 6 June



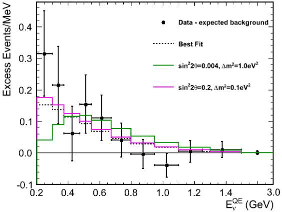
fit of low-energy excesses is excluded by reactor bound!

Total Excess: 240.3 +/- 34.5 +/- 52.6



* Simultaneous fit (E>200 MeV) with fully-correlated systematic to entire MB neutrino and anti-neutrino data

combined	E > 200 MeV	E > 475 MeV
χ^2(null)	42.53	12.87
Prob(null)	0.1%	35.8%
χ^2(bf)	24.72	10.67
Prob(bf)	6.7%	35.8%

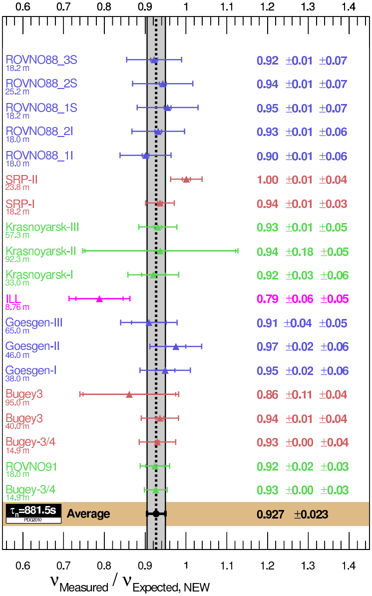


Neutrino energy reconstruction problem?

[Martini, Ericson, Chanfray, arXiv:1202.4745]

Reactor Electron Antineutrino Anomaly

[Mention et al, PRD 83 (2011) 073006; update in White Paper, arXiv:1204.5379]



new reactor $\bar{\nu}_e$ fluxes

[Mueller et al, PRC 83 (2011) 054615]

[Huber, PRC 84 (2011) 024617]

Detection: $\sigma(\bar{\nu}_e + p \to n + e^+) \propto \tau_n^{-1}$

PDG	neutron lifetime τ_n
1995	887.0 ± 2.0 sec
1998	886.7 ± 1.9 sec
2002	885.7 ± 0.8 sec
2011	881.5 ± 1.5 sec
2012	880.1 ± 1.1 sec

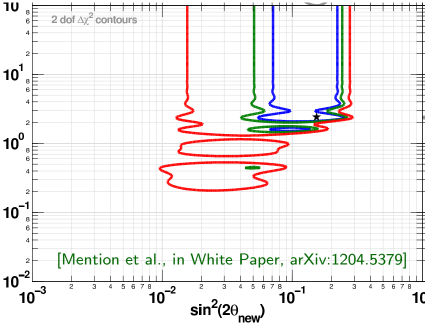
change of predicted event rates

^{235}U	^{238}U	^{239}Pu	^{241}Pu
$+3.7\%$	$+9.8\%$	$+4.2\%$	$+4.7\%$

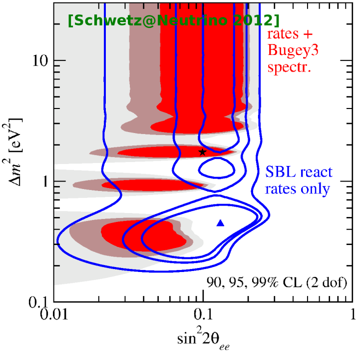
Reactor $\bar{\nu}_e$ Disappearance

$$P_{\nu_e \to \nu_e} = 1 - \sin^2 2\vartheta_{ee} \sin^2\left(\frac{\Delta m^2_{41} L}{4E}\right)$$

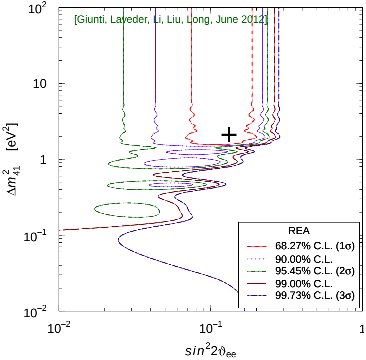
$$\sin^2 2\vartheta_{ee} = 4|U_{e4}|^2 \left(1 - |U_{e4}|^2\right)$$



[Mention et al., in White Paper, arXiv:1204.5379]



[Schwetz@Neutrino 2012]



[Giunti, Laveder, Li, Liu, Long, June 2012]

Gallium Anomaly

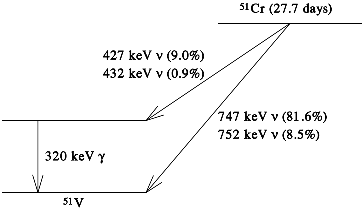
Gallium Radioactive Source Experiments

Tests of the solar neutrino detectors GALLEX (Cr1, Cr2) and SAGE (Cr, Ar)

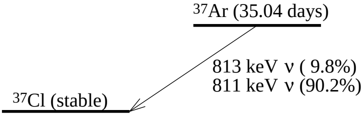
Detection Process: $\nu_e + {}^{71}\text{Ga} \to {}^{71}\text{Ge} + e^-$

ν_e Sources: $e^- + {}^{51}\text{Cr} \to {}^{51}\text{V} + \nu_e$ $e^- + {}^{37}\text{Ar} \to {}^{37}\text{Cl} + \nu_e$

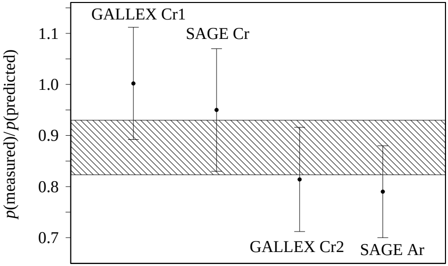
	^{51}Cr				^{37}Ar	
E [keV]	747	752	427	432	811	813
B.R.	0.8163	0.0849	0.0895	0.0093	0.902	0.098



[SAGE, PRC 59 (1999) 2246, hep-ph/9803418]



[SAGE, PRC 73 (2006) 045805, nucl-ex/0512041]



[SAGE, PRC 73 (2006) 045805, nucl-ex/0512041]

$$
\begin{aligned}
R_{\mathrm{B}}^{\text{GALLEX-Cr1}} &= 0.953 \pm 0.11 \\
R_{\mathrm{B}}^{\text{GALLEX-Cr2}} &= 0.812^{+0.10}_{-0.11} \\
R_{\mathrm{B}}^{\text{SAGE-Cr}} &= 0.95 \pm 0.12 \\
R_{\mathrm{B}}^{\text{SAGE-Ar}} &= 0.791^{+0.084}_{-0.078}
\end{aligned}
$$

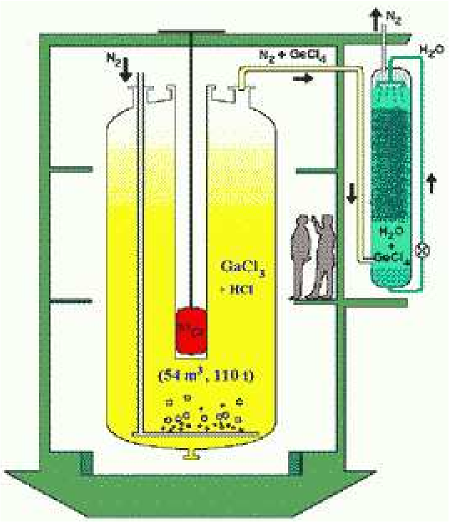
$$\boxed{R_{\mathrm{B}}^{\mathrm{Ga}} = 0.86 \pm 0.05}$$

Bahcall cross section without uncertainty

[Bahcall, PRC 56 (1997) 3391, hep-ph/9710491]

$$\langle L \rangle_{\mathrm{GALLEX}} = 1.9\,\mathrm{m}$$

$$\langle L \rangle_{\mathrm{SAGE}} = 0.6\,\mathrm{m}$$

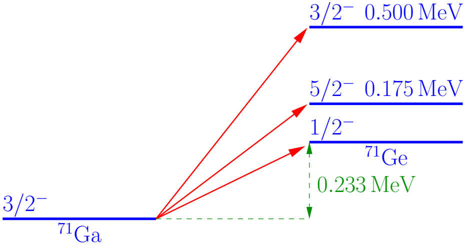


[GALLEX]

- Deficit could be due to overestimate of

$$\sigma(\nu_e + {}^{71}\text{Ga} \to {}^{71}\text{Ge} + e^-)$$

- Calculation: Bahcall, PRC 56 (1997) 3391



- $\sigma_{\text{G.S.}}$ from $T_{1/2}({}^{71}\text{Ge}) = 11.43 \pm 0.03\,\text{days}$ [Hampel, Remsberg, PRC 31 (1985) 666]

$$\sigma_{\text{G.S.}}({}^{51}\text{Cr}) = 55.3 \times 10^{-46}\,\text{cm}^2\,(1 \pm 0.004)_{3\sigma}$$

- $\sigma({}^{51}\text{Cr}) = \sigma_{\text{G.S.}}({}^{51}\text{Cr}) \left(1 + 0.669\,\frac{\text{BGT}_{175}}{\text{BGT}_{\text{G.S.}}} + 0.220\,\frac{\text{BGT}_{500}}{\text{BGT}_{\text{G.S.}}}\right)$

- Contribution of Excited States only 5%!

		$\frac{\mathrm{BGT}_{175}}{\mathrm{BGT}_{\mathrm{G.S.}}}$	$\frac{\mathrm{BGT}_{500}}{\mathrm{BGT}_{\mathrm{G.S.}}}$
Krofcheck et al. PRL 55 (1985) 1051	$^{71}\mathrm{Ga}(p,n)^{71}\mathrm{Ge}$	< 0.056	0.13 ± 0.02
Haxton PLB 431 (1998) 110	Shell Model	0.19 ± 0.18	
Frekers et al. PLB 706 (2011) 134	$^{71}\mathrm{Ga}(^{3}\mathrm{He},{}^{3}\mathrm{H})^{71}\mathrm{Ge}$	0.039 ± 0.030	0.202 ± 0.016

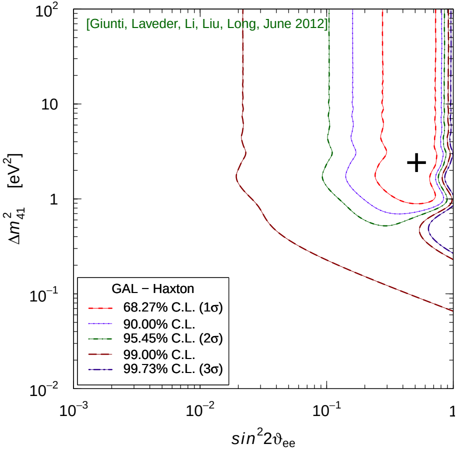
- Haxton: [Haxton, PLB 431 (1998) 110]

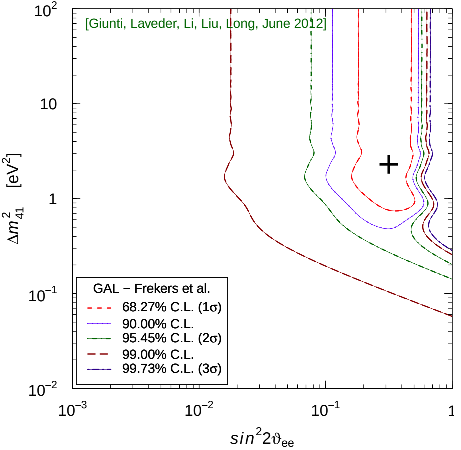
  “a sophisticated shell model calculation is performed … for the transition to the first excited state in $^{71}\mathrm{Ge}$. The calculation predicts destructive interference between the (p,n) spin and spin-tensor matrix elements”

- 2.7σ discrepancy of $\mathrm{BGT}_{500}/\mathrm{BGT}_{\mathrm{G.S.}}$ measurements

- Anyhow, new $^{71}\mathrm{Ga}(^{3}\mathrm{He},{}^{3}\mathrm{H})^{71}\mathrm{Ge}$ data support Gallium Anomaly!

Gallium ν_e Disappearance

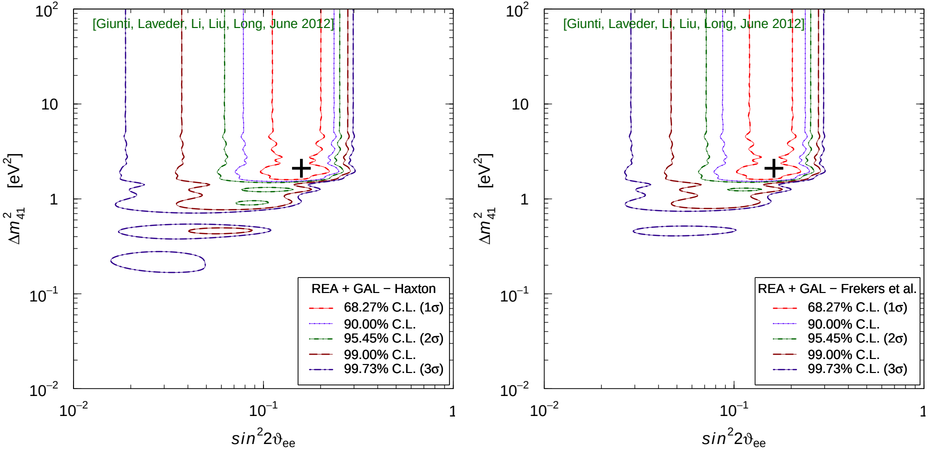




No Osc.	$\chi^2_{\rm min}$	14.9
	NDF	4
	GoF	0.5 %
3+1	$\chi^2_{\rm min}$	4.7
	NDF	2
	GoF	9.5 %
	Δm^2_{41}[eV^2]	2.24
	$\sin^2 2\vartheta_{ee}$	0.51

No Osc.	$\chi^2_{\rm min}$	18.2
	NDF	4
	GoF	0.1 %
3+1	$\chi^2_{\rm min}$	7.9
	NDF	2
	GoF	1.9 %
	Δm^2_{41}[eV^2]	2.14
	$\sin^2 2\vartheta_{ee}$	0.32

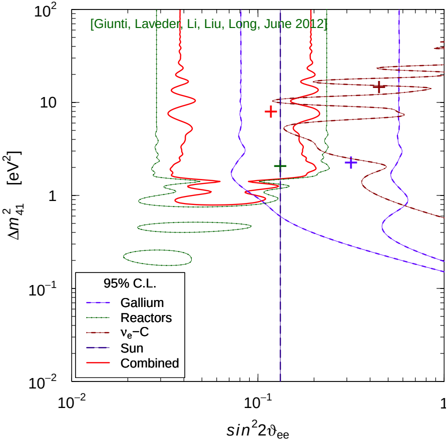
Reactor $\bar{\nu}_e$ and Gallium ν_e Disappearance

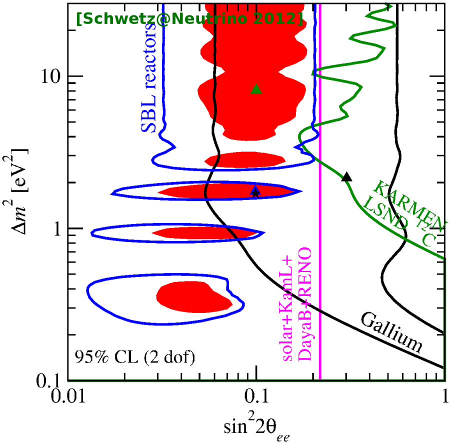


No Osc.	χ^2_{min}	45.6
	NDF	42
	GoF	32.3 %
3+1	χ^2_{min}	30.8
	NDF	40
	GoF	85 %
	Δm^2_{41} [eV^2]	1.95
	$\sin^2 2\vartheta_{ee}$	0.16

No Osc.	χ^2_{min}	48.9
	NDF	42
	GoF	21.5 %
3+1	χ^2_{min}	32.2
	NDF	40
	GoF	80 %
	Δm^2_{41} [eV^2]	1.95
	$\sin^2 2\vartheta_{ee}$	0.16

Global ν_e and $\bar{\nu}_e$ Disappearance





KARMEN + LSND

$\nu_e + {}^{12}C \to {}^{12}N_{g.s.} + e^-$

[Conrad, Shaevitz, PRD 85 (2012) 013017]

[Giunti, Laveder, PLB 706 (2011) 200]

SUN&KamLAND + ϑ_{13}

[Giunti, Li, PRD 80 (2009) 113007]

[Palazzo, PRD 83 (2011) 113013]

[Palazzo, PRD 85 (2012) 077301]

SUN&KamLAND + ϑ_{13} bound on $|U_{e4}|^2$

[Giunti, Li, PRD 80 (2009) 113007; Palazzo, PRD 83 (2011) 113013, PRD 85 (2012) 077301]

3+1 with simplifying assumptions: $U_{\mu 4} = U_{\tau 4} = 0$, no CP violation

$$U_{e1} = c_{12}c_{13}c_{14} \quad U_{e2} = s_{12}c_{13}c_{14} \quad U_{e3} = s_{13}c_{14} \quad U_{e4} = s_{14}$$

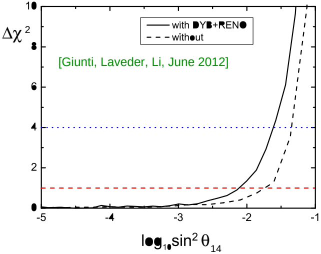
$$U_{s1} = -c_{12}c_{13}s_{14} \quad U_{s2} = -s_{12}c_{13}s_{14} \quad U_{s3} = -s_{13}s_{14} \quad U_{s4} = c_{14}$$

$$P_{\nu_e \to \nu_e} = c_{13}^4 c_{14}^4 P_{\nu_e \to \nu_e}^{2\nu} + s_{13}^4 c_{14}^4 + s_{14}^4$$

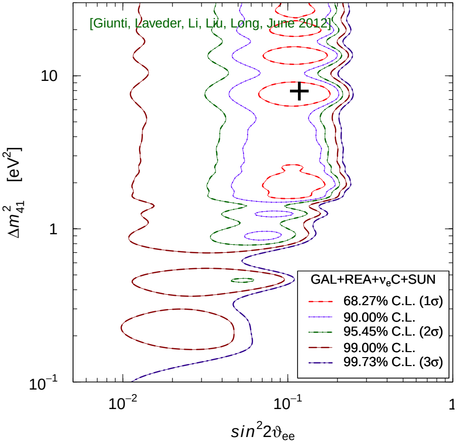
$$P_{\nu_e \to \nu_s} = c_{14}^2 s_{14}^2 \left(c_{13}^4 P_{\nu_e \to \nu_s}^{2\nu} + s_{13}^4 + 1 \right)$$

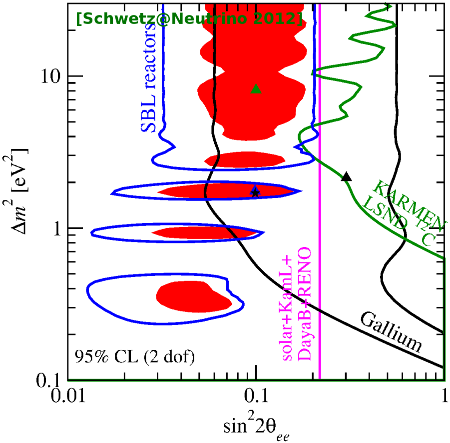
$$V = c_{13}^2 c_{14}^2 V_{\text{CC}} - c_{13}^2 s_{14}^2 V_{\text{NC}} = \left(|U_{e1}|^2 + |U_{e2}|^2\right) V_{\text{CC}} - \left(|U_{s1}|^2 + |U_{s2}|^2\right) V_{\text{NC}}$$

Fit with $U_{\mu 4}$ and $U_{\tau 4}$ free:



Global ν_e and $\bar{\nu}_e$ Disappearance



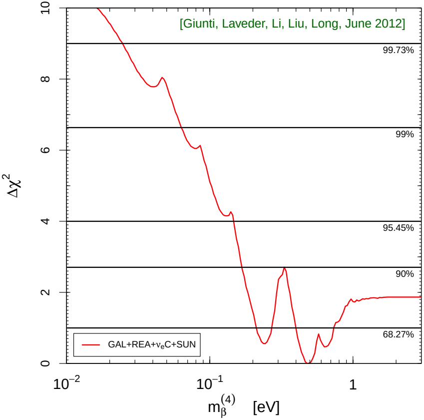


No Osc.	$\chi^2_{\rm min}$	57.1
	NDF	53
	GoF	32.6 %
3+1	$\chi^2_{\rm min}$	46.0
	NDF	51
	GoF	67 %
	Δm^2_{41}[eV^2]	7.59
	$\sin^2 2\vartheta_{ee}$	0.12

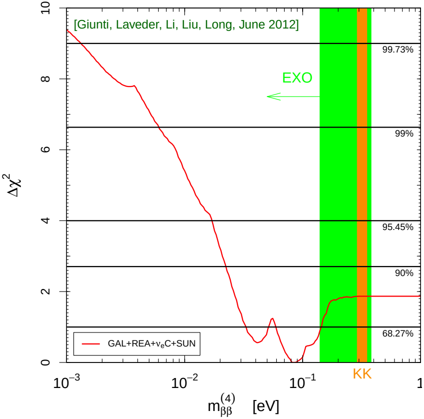
No Osc.	$\chi^2_{\rm min}$	318.4
	NDF	331
	GoF	68%
3+1	$\chi^2_{\rm min}$	306.0
	NDF	329
	GoF	80%
	Δm^2_{41}[eV^2]	1.71
	$\sin^2 2\vartheta_{ee}$	0.099

Testable Implications

β Decay



$(\beta\beta)_{0\nu}$ Decay



$$m_\beta = \sqrt{\sum_k |U_{ek}|^2 m_k^2}$$

$$m_\beta^{(4)} = |U_{e4}|\sqrt{\Delta m_{41}^2}$$

$$m_{\beta\beta} = \left|\sum_k U_{ek}^2\, m_k\right|$$

$$m_{\beta\beta}^{(4)} = |U_{e4}|^2\sqrt{\Delta m_{41}^2}$$

Global 3+1 Fit: Disappearance Constraints

- ν_e disappearance experiments:

$$\sin^2 2\vartheta_{ee} = 4|U_{e4}|^2 \left(1 - |U_{e4}|^2\right) \simeq 4|U_{e4}|^2$$

- ν_μ disappearance experiments:

$$\sin^2 2\vartheta_{\mu\mu} = 4|U_{\mu4}|^2 \left(1 - |U_{\mu4}|^2\right) \simeq 4|U_{\mu4}|^2$$

- $\nu_\mu \to \nu_e$ experiments:

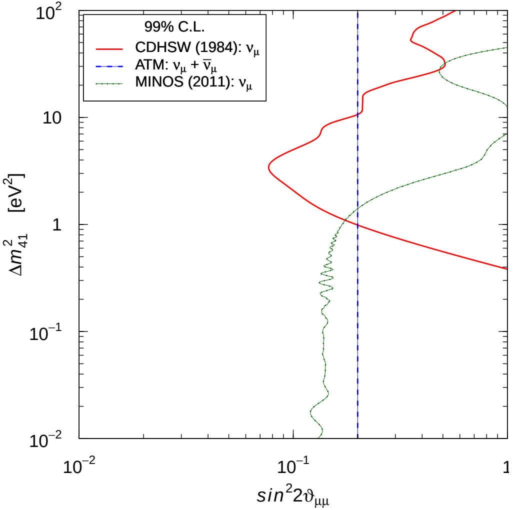
$$\sin^2 2\vartheta_{e\mu} = 4|U_{e4}|^2|U_{\mu4}|^2 \simeq \frac{1}{4} \sin^2 2\vartheta_{ee} \sin^2 2\vartheta_{\mu\mu}$$

- Upper bounds on $\sin^2 2\vartheta_{ee}$ and $\sin^2 2\vartheta_{\mu\mu}$ $\Longrightarrow$ strong limit on $\sin^2 2\vartheta_{e\mu}$

[Okada, Yasuda, Int. J. Mod. Phys. A12 (1997) 3669-3694]

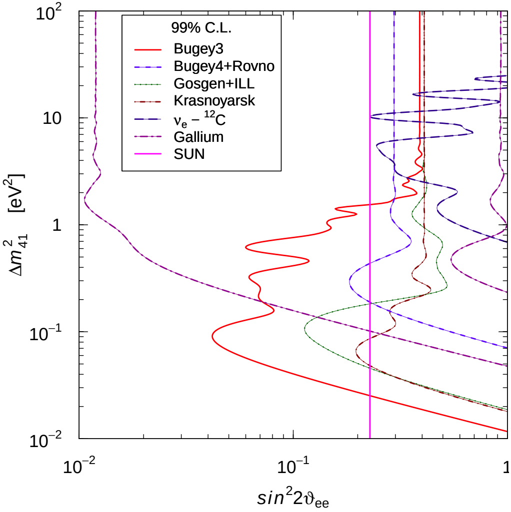
[Bilenky, Giunti, Grimus, Eur. Phys. J. C1 (1998) 247]

ν_μ and $\bar{\nu}_\mu$ Disappearance



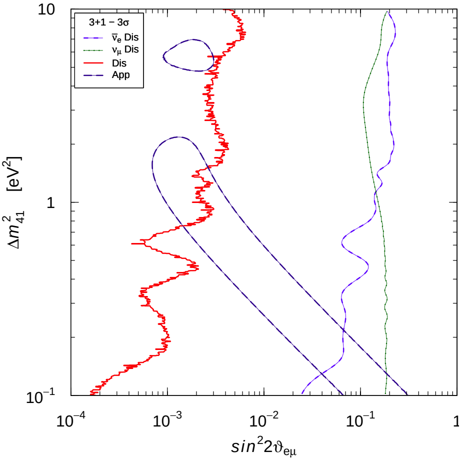
- ATM constraint on $|U_{\mu 4}|^2$
  [Maltoni, Schwetz, PRD 76 (2007) 093005]
- MINOS constraint on $|U_{\mu 4}|^2$
  [Giunti, Laveder, PRD 84 (2011) 093006]

ν_e and $\bar{\nu}_e$ Disappearance



- New Reactor $\bar{\nu}_e$ Fluxes
  [Mueller et al., PRC 83 (2011) 054615]
  [Mention et al., PRD 83 (2011) 073006]
  [Huber, PRC 84 (2011) 024617]
- KARMEN + LSND
  $\nu_e + {}^{12}\mathrm{C} \to {}^{12}\mathrm{N}_{\mathrm{g.s.}} + e^-$
  [Conrad, Shaevitz, PRD 85 (2012) 013017]
  [Giunti, Laveder, PLB 706 (2011) 200]
- SUN&KamLAND + ϑ_{13}
  [Giunti, Li, PRD 80 (2009) 113007]
  [Palazzo, PRD 83 (2011) 113013]
  [Palazzo, PRD 85 (2012) 077301]
  [Giunti, Laveder, Li, Liu, Long, in preparation]

3+1 with SUN&KamLAND+ϑ_{13}

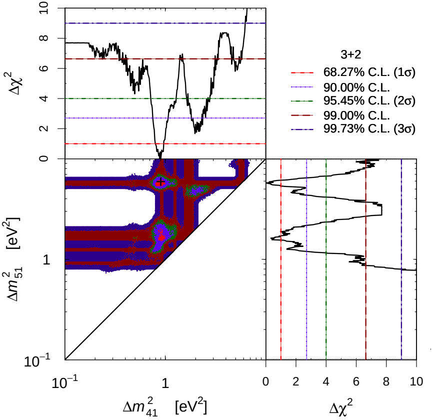


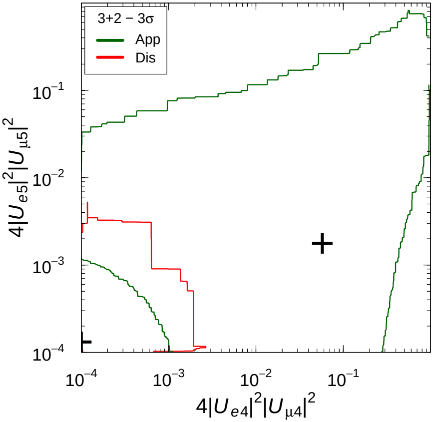
GoF = 38%

PGoF = 0.3%

- 3+1 & 3+2: Appearance-Disappearance tension
- Tension reduced in 3+1+NSI [Akhmedov, Schwetz, JHEP 10 (2010) 115]
- No tension in 3+1+CPTV [Barger, Marfatia, Whisnant, PLB 576 (2003) 303] [Giunti, Laveder, PRD 82 (2010) 093016, PRD 83 (2011) 053006]

3+2 with SUN&KamLAND+ϑ_{13}

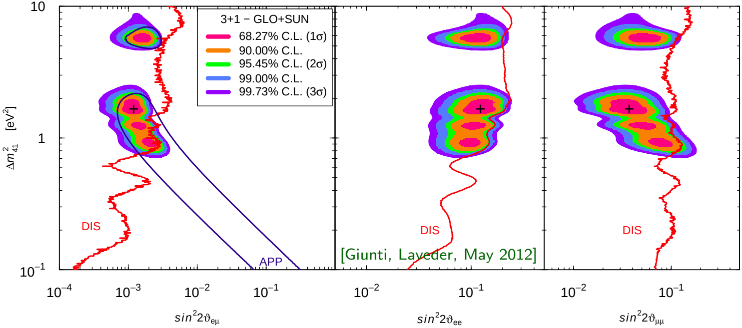




GoF = 46%   PGoF = 0.29%

- 3+2 is preferred to 3+1 only if there is CP-violating MiniBooNE neutrino-antineutrino difference
- almost disappeared with 2012 MiniBooNE antineutrino data

3+1 Global Fit



[Giunti, Laveder, May 2012]

- MiniBooNE 2011 data ($E > 475\,\mathrm{MeV}$)
- More time is needed to fit MiniBooNE 2012 data

Cosmology

- N_s = number of thermalized sterile neutrinos (not necessarily integer)
- CMB and LSS in ΛCDM: $N_s = 1.3 \pm 0.9$ $m_s < 0.66\,\mathrm{eV}$ (95% C.L.)

  [Hamann, Hannestad, Raffelt, Tamborra, Wong, PRL 105 (2010) 181301]

  $N_s = 1.61 \pm 0.92$ $m_s < 0.70\,\mathrm{eV}$ (95% C.L.)

  [Giusarma, Corsi, Archidiacono, de Putter, Melchiorri, Mena, Pandolfi, PRD 83 (2011) 115023]

  $N_s = 1.12^{+0.86}_{-0.74}$ (95% C.L.) [Archidiacono, Calabrese, Melchiorri, PRD 84 (2011) 123008]

- BBN: $\begin{cases} N_s = 0.22 \pm 0.59 \\ N_s = 0.64^{+0.40}_{-0.35} \\ N_s \leq 1 \text{ at 95\% C.L.} \end{cases}$

  [Cyburt, Fields, Olive, Skillman, AP 23 (2005) 313]

  [Izotov, Thuan, ApJL 710 (2010) L67]

  [Mangano, Serpico, PLB 701 (2011) 296]

- CMB+LSS+BBN: $N_s = 0.85^{+0.39}_{-0.56}$ (95% C.L.)

  [Hamann, Hannestad, Raffelt, Wong, JCAP 1109 (2011) 034]

- Standard ΛCDM: 3+1 allowed, 3+2 disfavored

Conclusions

- Short-baseline neutrino oscillation anomalies $\Longrightarrow$ sterile neutrinos
- Short-Baseline $\bar{\nu}_\mu \to \bar{\nu}_e$ Signal is not feeling well:
  - MiniBooNE 2011 antineutrino data were more similar to neutrino data than those of 2010 (LSND signal diminished and low-energy anomaly appeared)
  - MiniBooNE 2012 antineutrino data are even more similar to neutrino data
  - Probably there is no CP violation $\Longrightarrow$ no need of 3+2
  - The decrease of MiniBooNE-LSND agreement is discouraging
  - Better experiments are needed to clarify situation
- Short-Baseline ν_e and $\bar{\nu}_e$ Disappearance is in good health:
  - Reactor $\bar{\nu}_e$ anomaly is alive and exciting
  - Gallium ν_e anomaly has been strengthened by new cross-section measurements
  - Many promising projects to test short-baseline ν_e and $\bar{\nu}_e$ disappearance in a few years with reactors and radioactive sources
  - Independent tests through effect of m_4 in β-decay and $(\beta\beta)_{0\nu}$-decay