

# Spreading width of doorway states and banded random matrices

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## Strength function

Strength function

$$S(E) = \sum_i |\langle 0|i \rangle|^2 \delta(E - \varepsilon_i) \quad (1)$$

where

$|0 \rangle$  doorway state

$\varepsilon_i$  and  $|i \rangle$  eigenvalues and eigenstates of the Hamiltonian  $H$

$S(E)$  displays a broad maximum  $\rightarrow E_0$

Important quantities:

$\Gamma \downarrow \rightarrow$  spreading width ( $\sim 4 \div 5$  MeV)

$d \rightarrow$  mean spacing of nuclear levels ( $\sim 10$  eV)

## Bohr-Mottelson

$$\overline{S(E)} = \frac{\Gamma^\downarrow / (2\pi)}{(E - E_0)^2 + (1/4)(\Gamma^\downarrow)^2} \quad (2)$$

with

$$\Gamma^\downarrow = 2\pi v^2 / d \quad (3)$$

Assumptions:

- constant level spacing  $d$  of the  $|i\rangle$
- coupling between background and  $|0\rangle \rightarrow V_\mu$  (Gaussian distributed with average 0 and variance  $v^2$ )

Our previous approach:

$$H = \begin{pmatrix} E_0 & V_\nu \\ V_\mu & \mathcal{H}_{\mu\nu} \end{pmatrix} \quad (4)$$

$H_{\mu\nu} \rightarrow$  GOE (real symmetric)

$|0\rangle$  at energy  $E_0$  coupled to  $N$  background states via  $V_\mu$  ( $\mu = 1..N$ ), real matrix elements

Results:

- spectrum of  $H_{\mu\nu} \rightarrow$  semicircle of radius  $2\lambda$
- at variance with Bohr-Mottelson, spectrum **is confined**
- solution obtained via Pastur equation if the spreading width is not negligible with respect to  $2\lambda$

Then  $\Gamma_{\text{eff}} > 2\pi v^2/d$ , the increase being proportional to  $\Gamma^\downarrow/\lambda$

Our new approach:

General background level distribution

$$H = \begin{pmatrix} E_0 & V_\nu \\ V_\mu & E_\mu \delta_{\mu\nu} \end{pmatrix} \quad (5)$$

$V_\mu$  Gaussian random variables with average 0 and variance  $v^2$

No assumption on the distribution of  $E_\mu$

Some details of the calculation

$$S(E) = -\frac{1}{\pi} \Im \left( \langle 0 | \frac{1}{E^+ - H} | 0 \rangle \right) \quad (6)$$

↓

$$S(E) = -\frac{1}{\pi} \Im \left( \frac{1}{E^+ - E_0 - \sum_{\mu} V_{\mu} (E^+ - E_{\mu})^{-1} V_{\mu}} \right) \quad (7)$$

↓

$$\bar{\Sigma}^V = v^2 \sum_{\mu} (E^+ - E_{\mu})^{-1} . \quad (8)$$

Finally

$$\overline{S(E)} = \frac{1}{2\pi} \frac{\Gamma^\downarrow}{(E - E_0 - \Delta)^2 + (1/4)(\Gamma^\downarrow)^2} \quad (9)$$

where

$$\begin{aligned} \Gamma^\downarrow &= 2\pi v^2 \rho(E) , \\ \Delta &= v^2 \int dE' \frac{\mathcal{P}}{E^+ - E'} \rho(E') \end{aligned} \quad (10)$$

and

$$\rho(E') = \sum_{\mu} \delta(E' - E_{\mu}) \quad (11)$$

The shift function is energy dependent

$$\rho(E) = \frac{N}{\pi\lambda} \sqrt{1 - \left(\frac{E}{2\lambda}\right)^2} \quad (\text{semicircle}),$$

$$\rho(E) = \frac{N}{\sqrt{2\pi\lambda}} \exp[-E^2/(2\lambda^2)] \quad (\text{Gaussian}) \quad (12)$$

↓

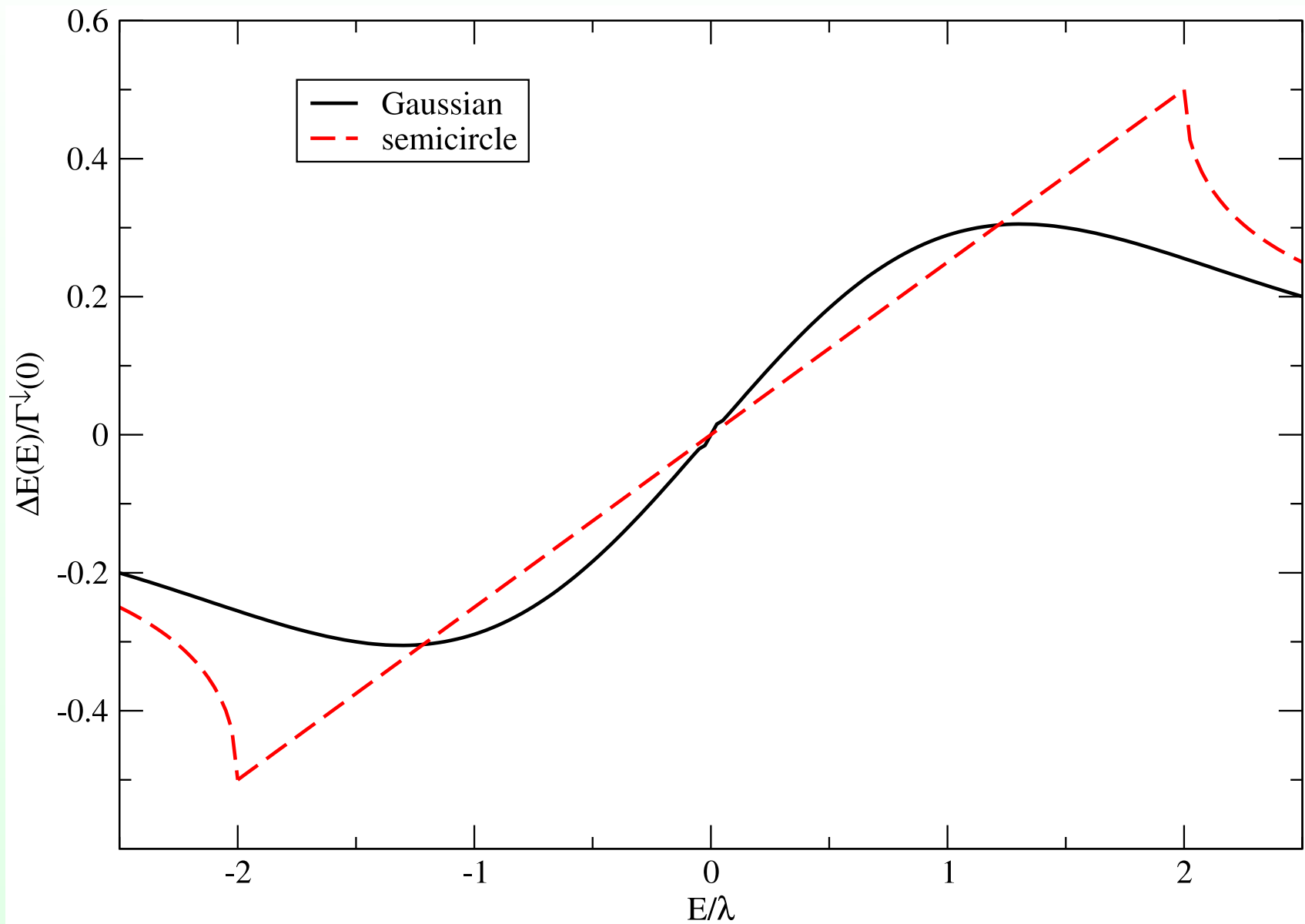
$$\Delta(E)/\Gamma^\downarrow(0) = \frac{1}{2\pi} \int_{-2}^{+2} dx' \frac{\mathcal{P}}{x - x'} \sqrt{1 - x'^2/4}$$

$$= \frac{x}{4} - [\theta(2 + x) - \theta(2 - x)] \frac{1}{2} \sqrt{\frac{x^2}{4} - 1}, \quad (\text{semicircle}),$$

$$\Delta(E)/\Gamma^\downarrow(0) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} dx' \frac{\mathcal{P}}{x - x'} \exp[-x'^2/2]$$

$$= \frac{1}{2} e^{-x^2/2} (-i) \operatorname{erf}\left(i \frac{x}{\sqrt{2}}\right) \quad (\text{Gaussian}) \quad (13)$$





## Numerical simulations

(A) GOE

because of the orthogonal invariance of GOE, individual matrix elements  $V_\mu$  are irrelevant, only

$$v^2 = \frac{1}{N} \sum_{\mu} V_{\mu}^2 \quad (14)$$

matters

$$\langle H_{\mu\nu} \rangle = 0 \text{ and } \overline{H_{\mu\nu}^2} = (1 + \delta_{\mu\nu}) \frac{\lambda^2}{N+1} \rightarrow \frac{1}{N} \text{Tr} \overline{H^2} = \lambda^2$$

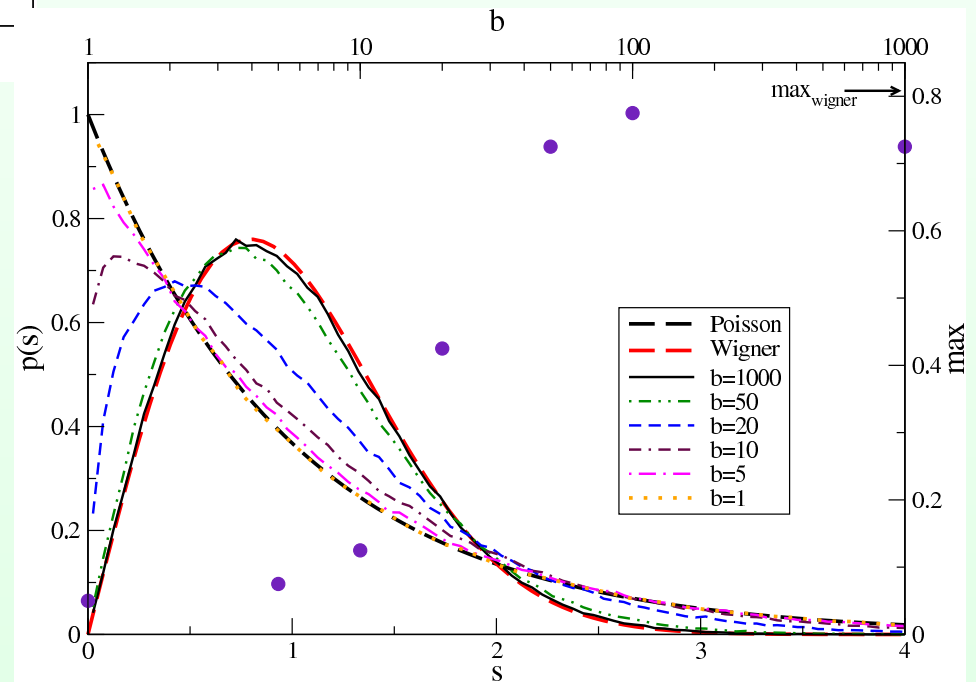
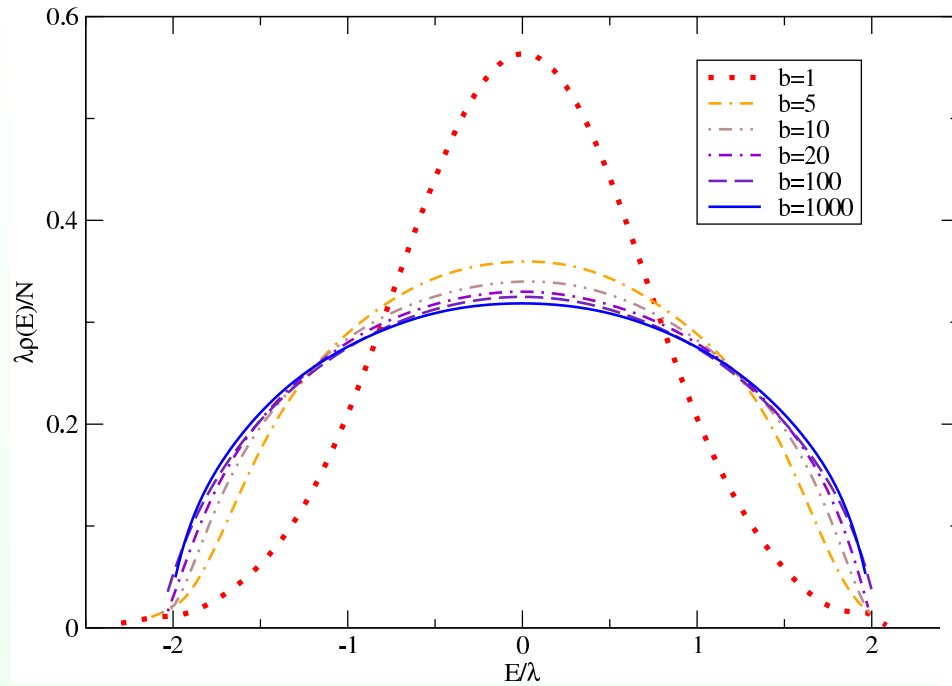
Parameters:  $v^2$ ,  $E_0$ ,  $N$  (=1000) (plus  $\lambda$ , which defines the energy scale and  $m$  (=500), the number of trials)

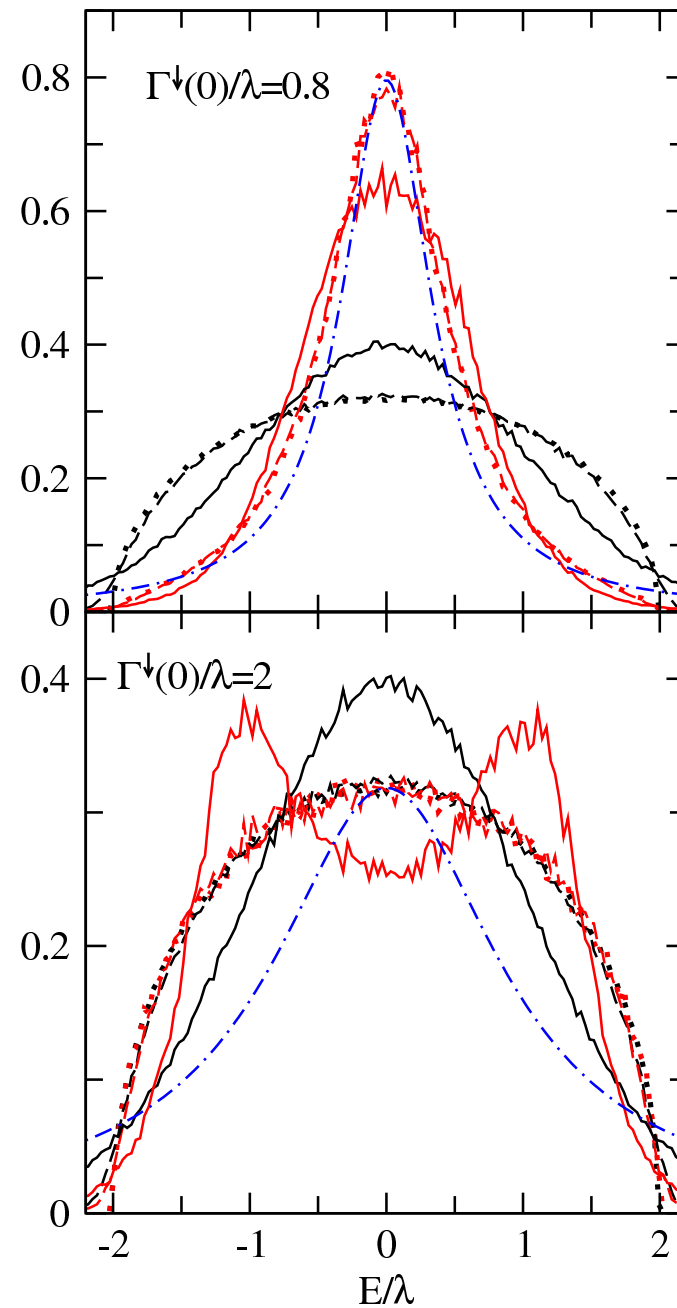
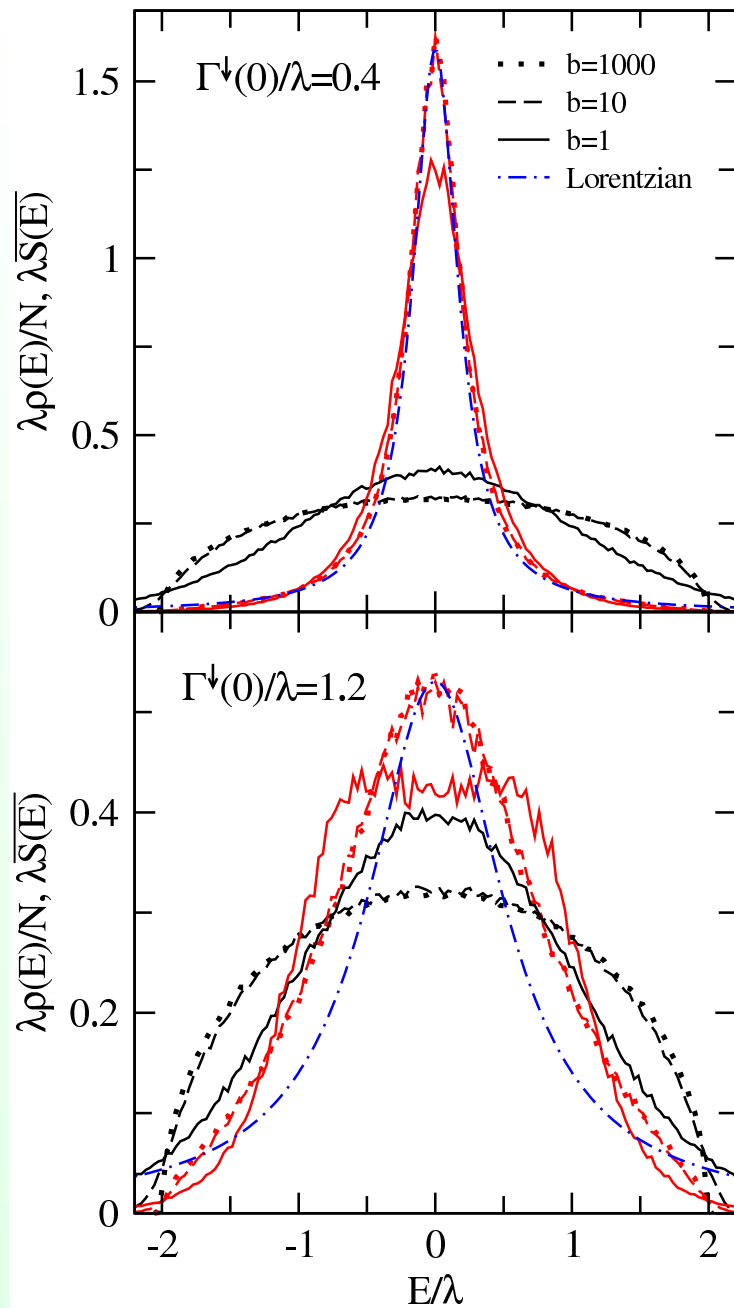
## Numerical simulations

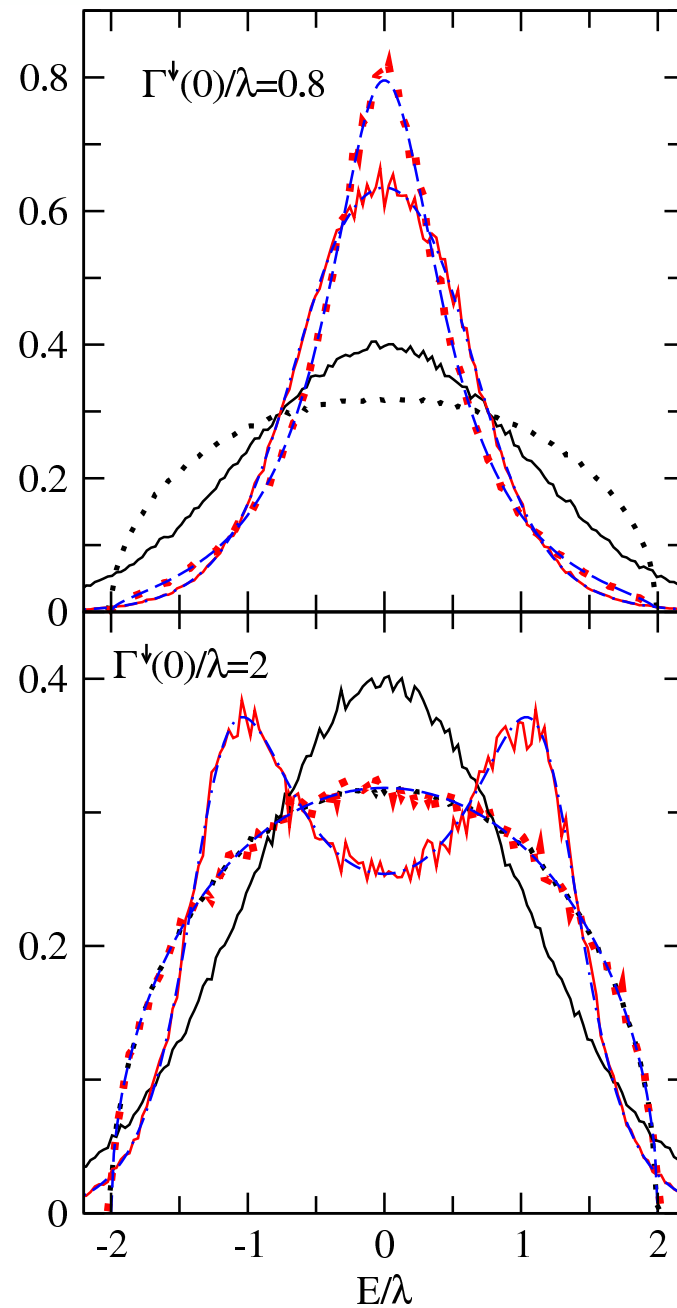
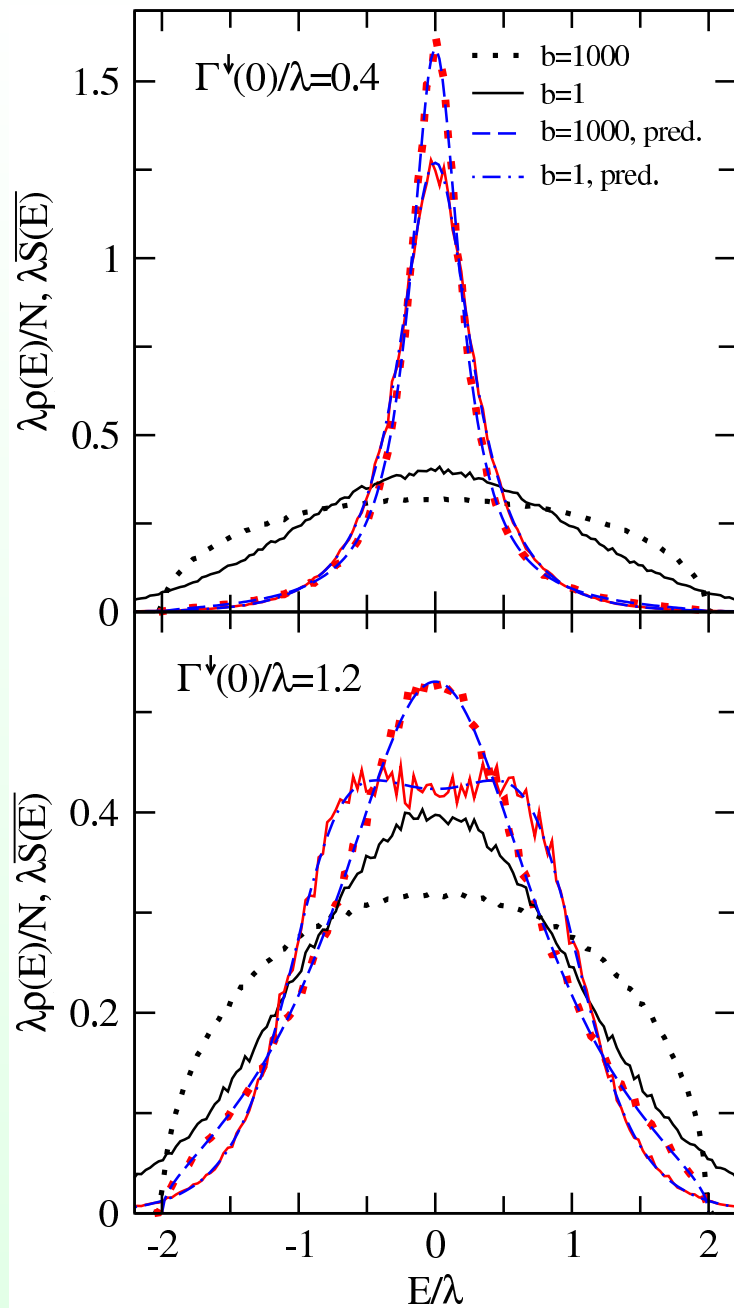
(B) banded matrices

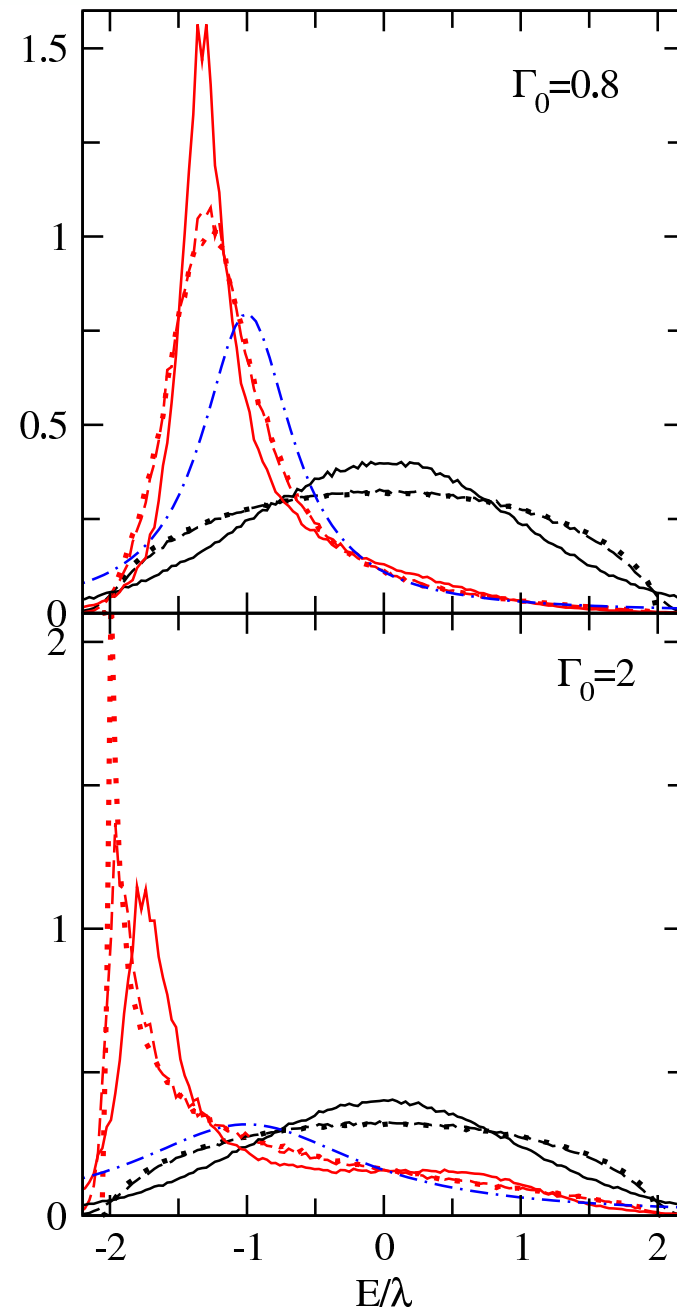
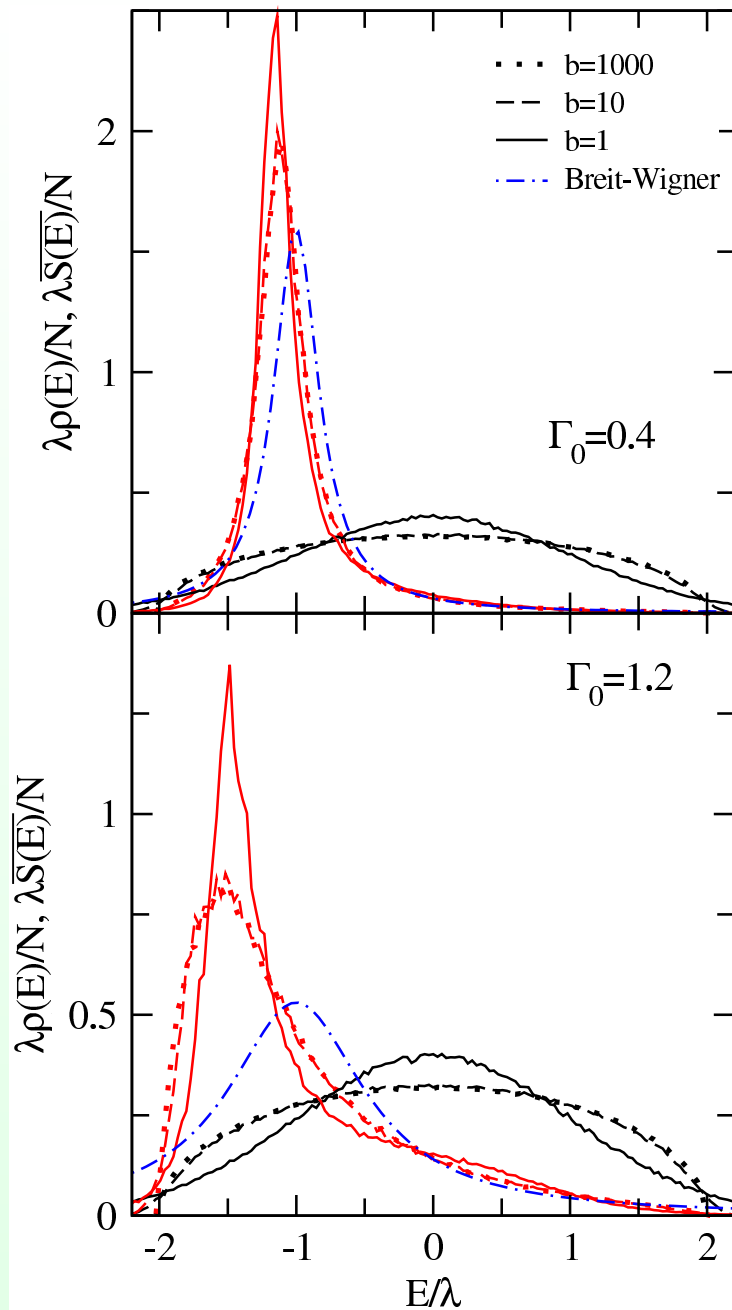
all matrix elements  $|\mu - \nu| \geq b$  vanish

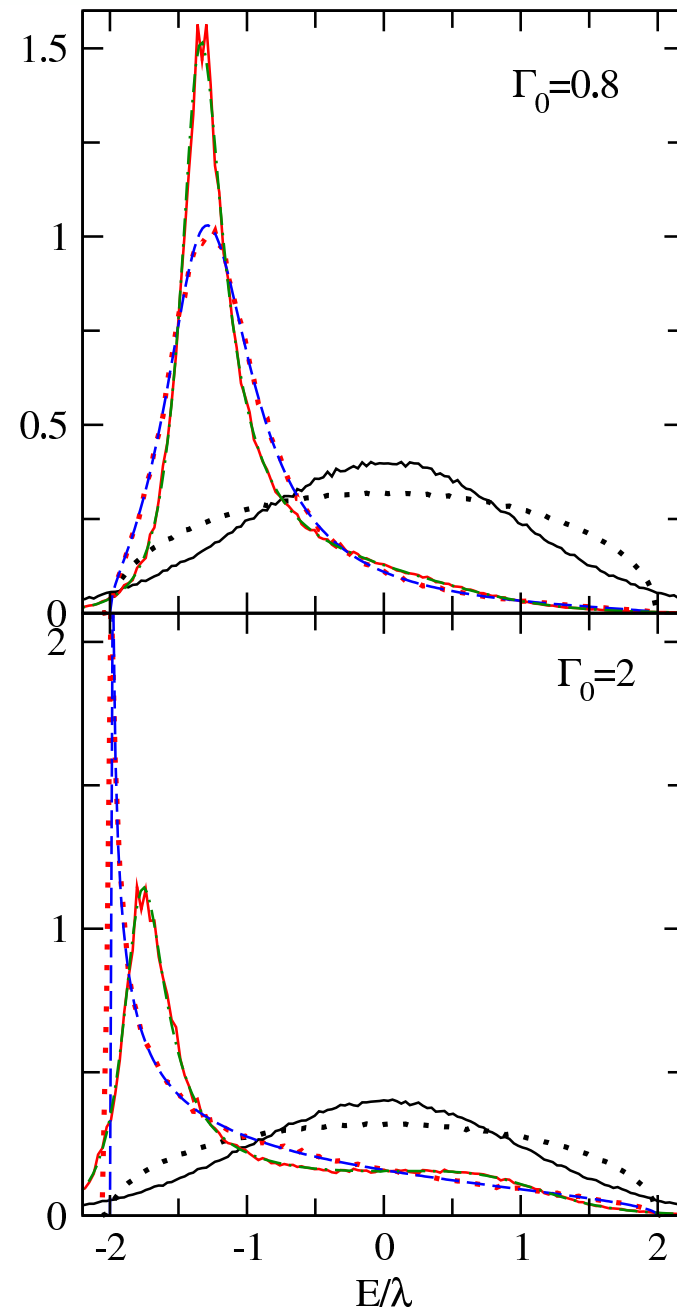
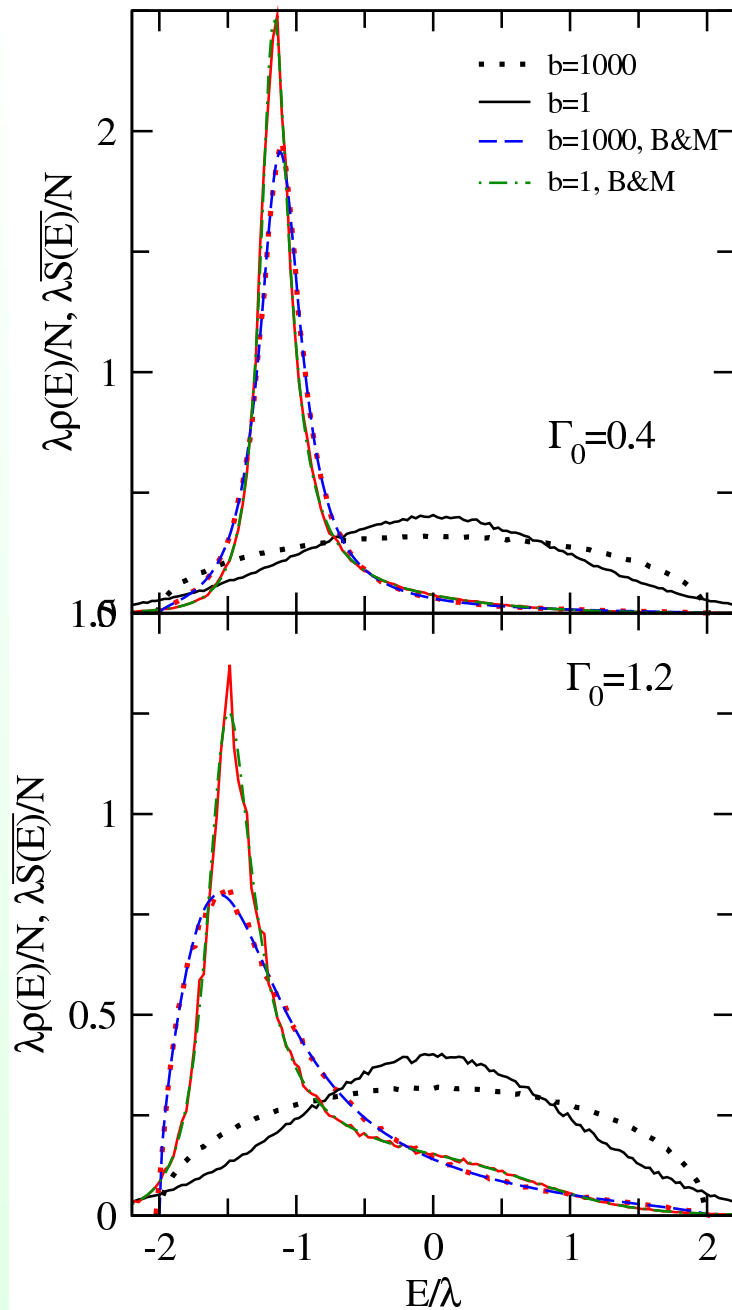
$$\overline{H_{\mu\nu}^2} = (1 + \delta_{\mu\nu})\beta^2 \leftarrow \frac{1}{N} \text{Tr} \overline{H^2} = \lambda^2$$













# Localization

Inverse participation ratio

$$P = \sum_i |\langle 0|i\rangle|^4$$

(15)

